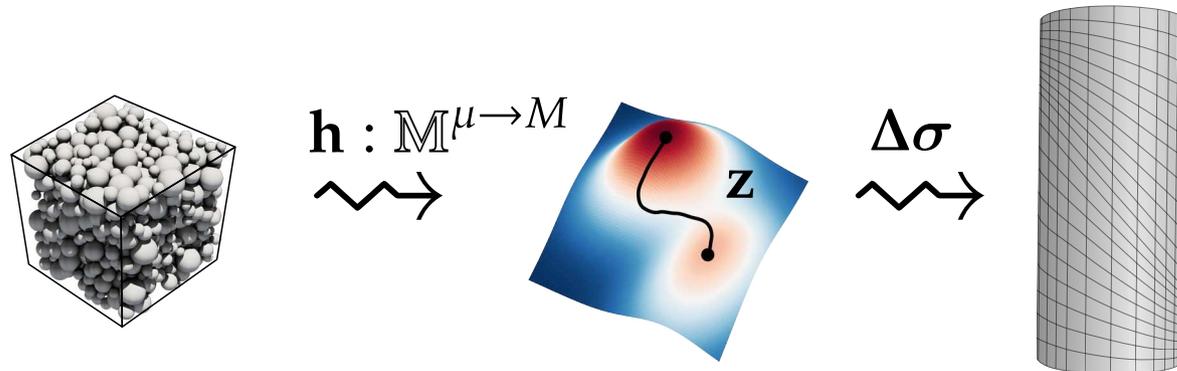


Thermodynamics- and data-driven scale bridging

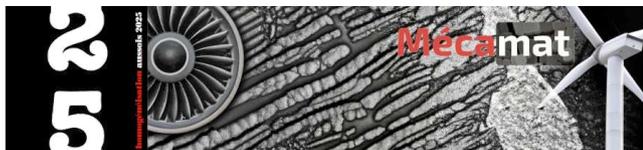


Filippo Masi

Univ. Grenoble Alpes, Inria, CNRS, Grenoble INP, LJK

Mécamat: Homogénéisation du comportement mécanique des matériaux hétérogènes

Aussois – Jan. 27-31, 2025



Scale-bridging

How to bridge the scales in inelastic, heterogeneous media

micro-

Macro-scale



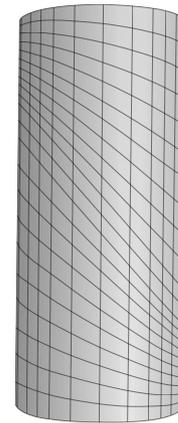
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media

micro-

Macro-scale

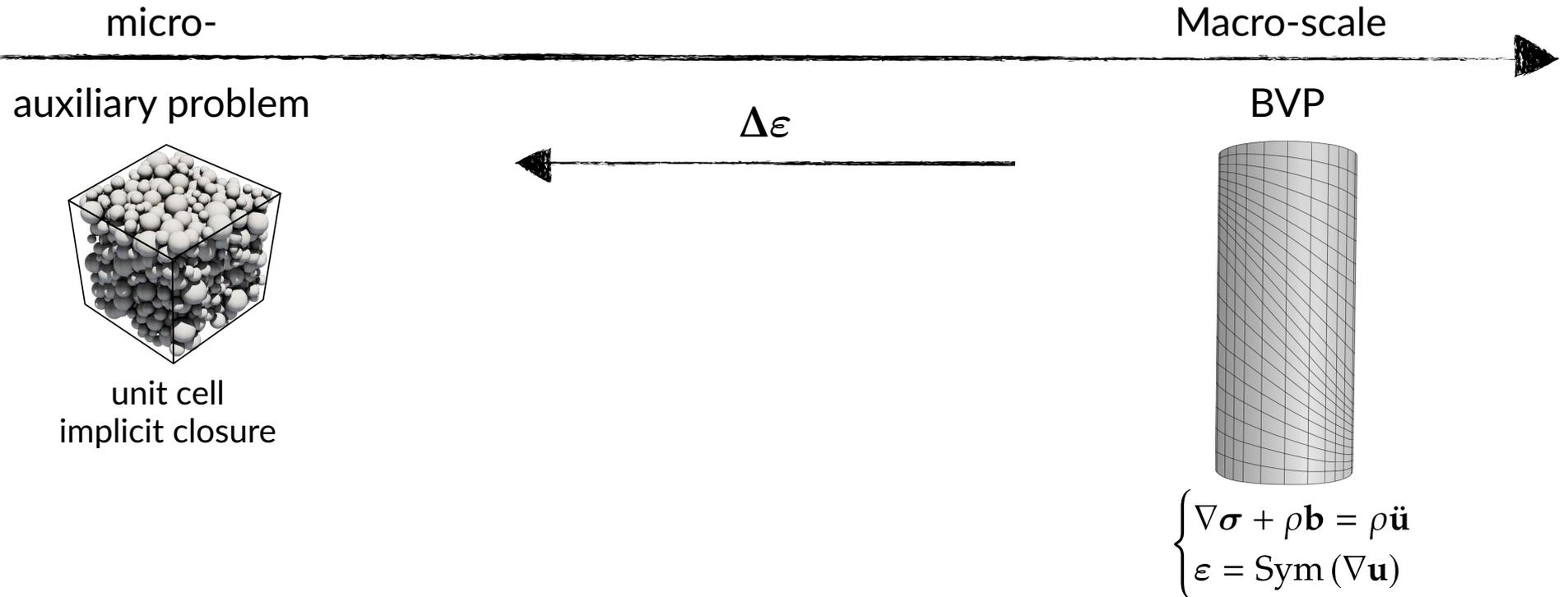
BVP



$$\begin{cases} \nabla \boldsymbol{\sigma} + \rho \mathbf{b} = \rho \ddot{\mathbf{u}} \\ \boldsymbol{\varepsilon} = \text{Sym}(\nabla \mathbf{u}) \end{cases}$$

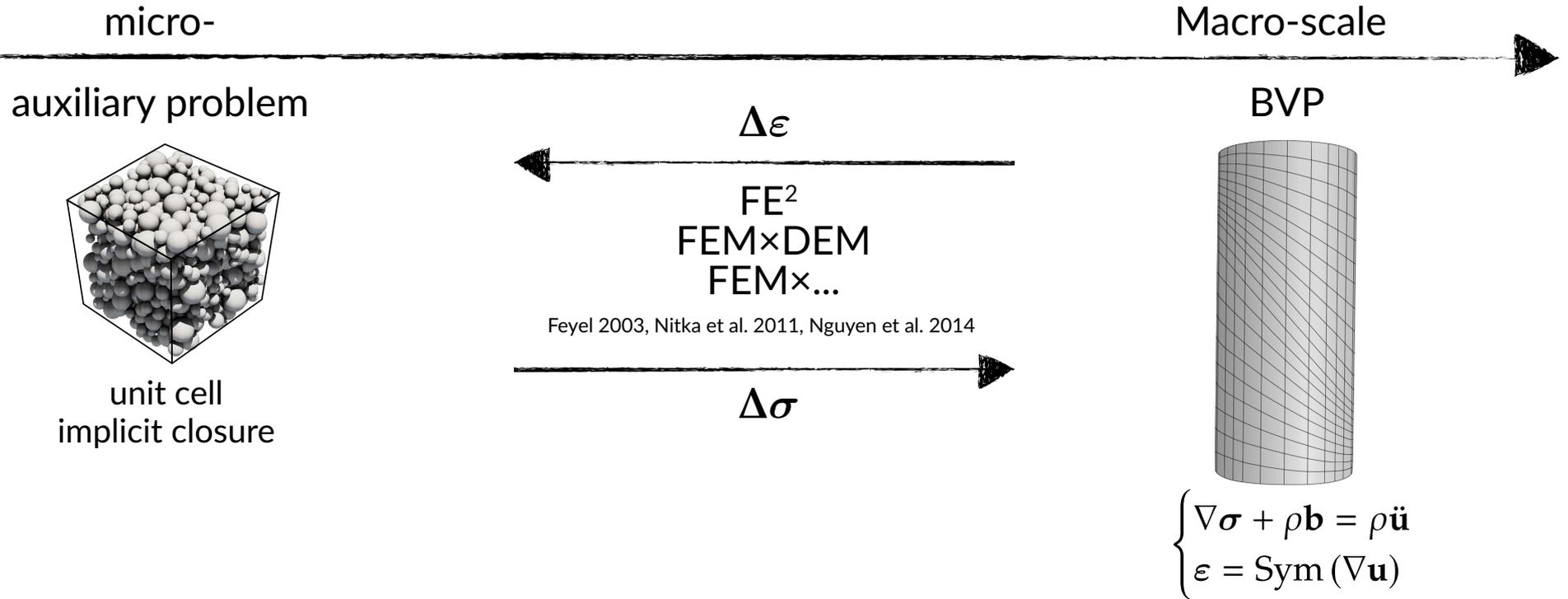
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media



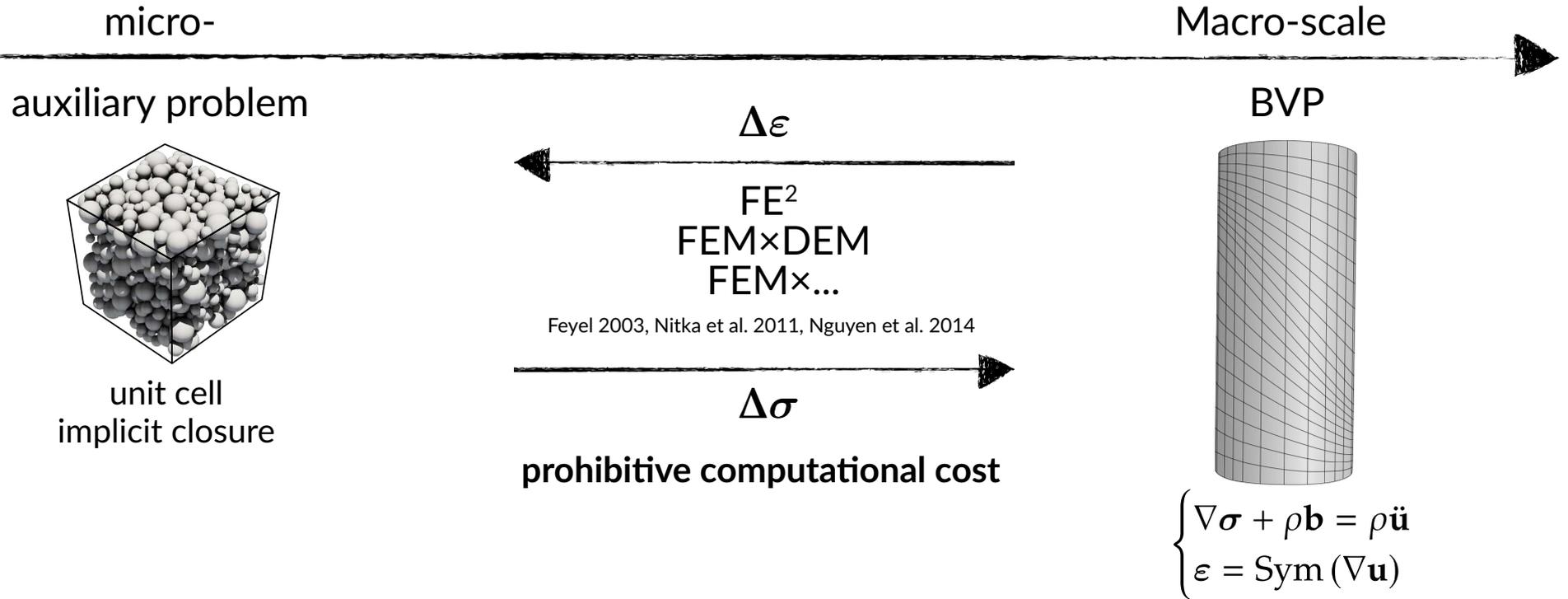
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media



Scale-bridging

How to bridge the scales in inelastic, heterogeneous media



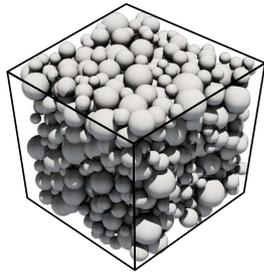
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media

micro-

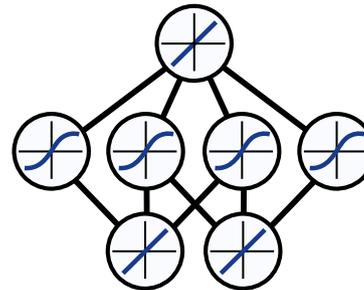
Macro-scale

auxiliary problem



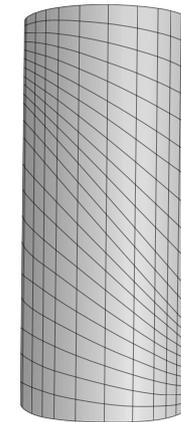
unit cell

ML (neural nets)



Ghaboussi et al. 1991,
Lefik & Schrefler 2003,
Mozaffar et al. 2019

BVP



$$\begin{cases} \nabla \boldsymbol{\sigma} + \rho \mathbf{b} = \rho \ddot{\mathbf{u}} \\ \boldsymbol{\varepsilon} = \text{Sym}(\nabla \mathbf{u}) \end{cases}$$

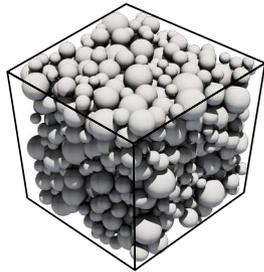
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media

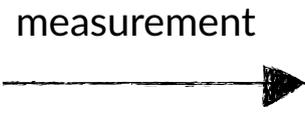
micro-

Macro-scale

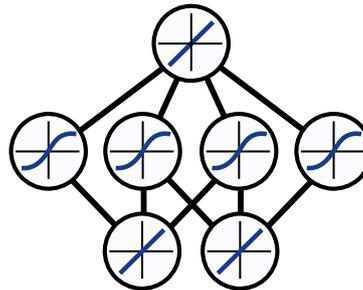
auxiliary problem



unit cell

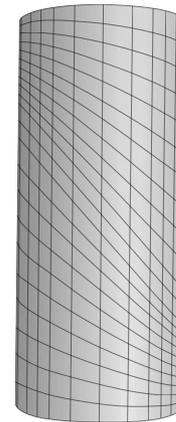


ML (neural nets)



Ghaboussi et al. 1991,
Lefik & Schrefler 2003,
Mozaffar et al. 2019

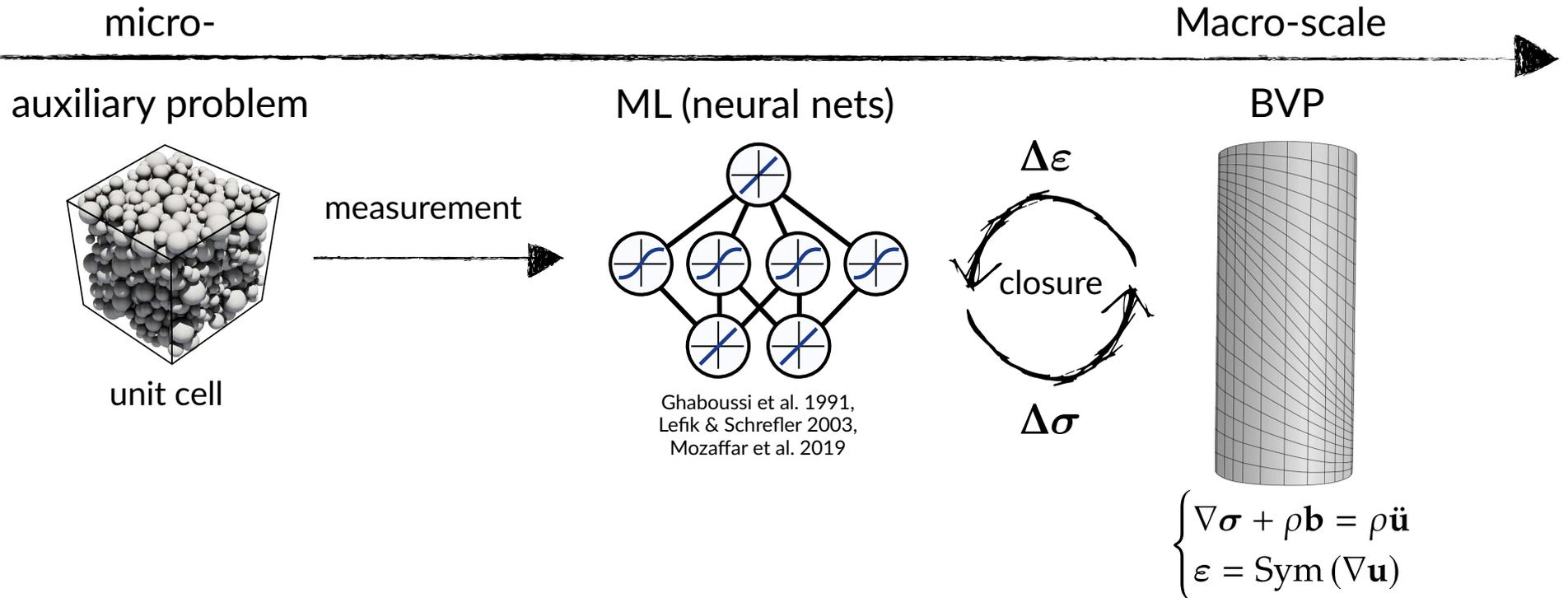
BVP



$$\begin{cases} \nabla \sigma + \rho \mathbf{b} = \rho \ddot{\mathbf{u}} \\ \varepsilon = \text{Sym}(\nabla \mathbf{u}) \end{cases}$$

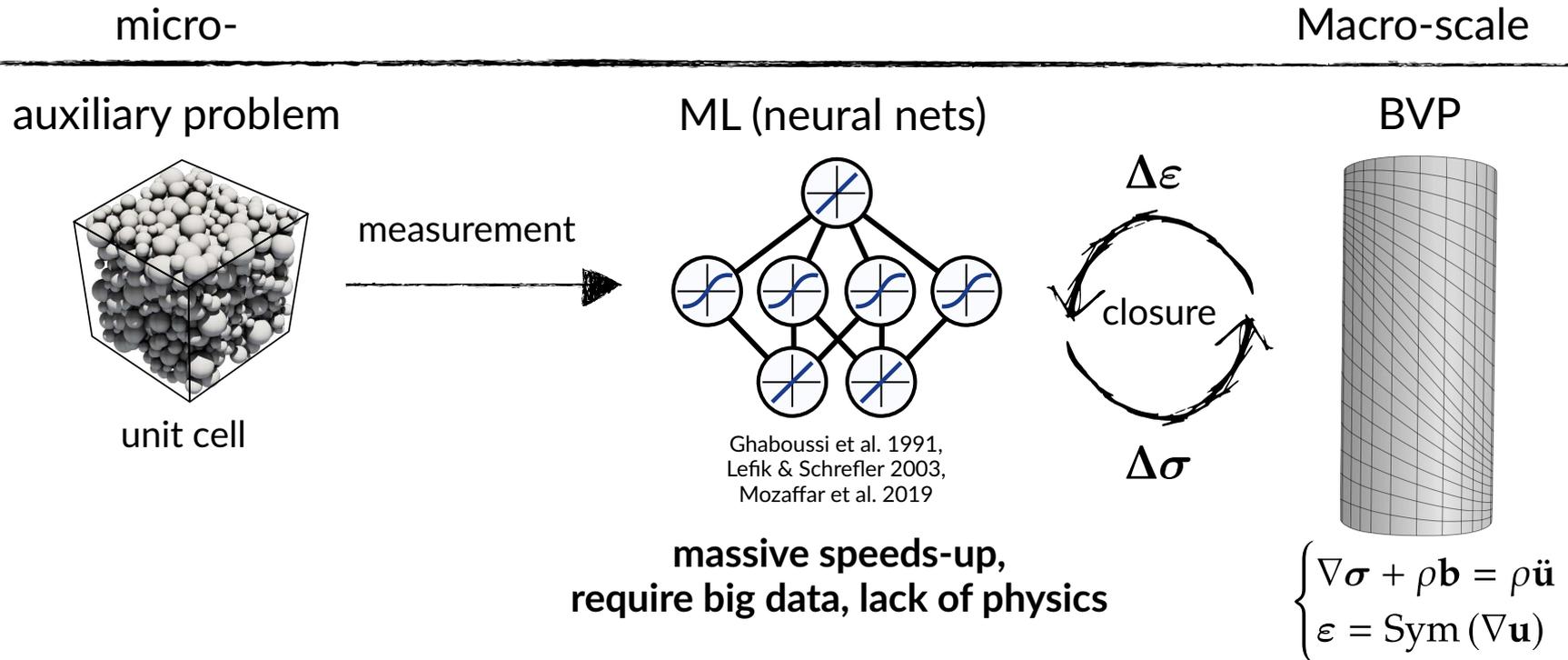
Scale-bridging

How to bridge the scales in inelastic, heterogeneous media



Scale-bridging

How to bridge the scales in inelastic, heterogeneous media



Scale-bridging

How to bridge the scales in inelastic, heterogeneous media

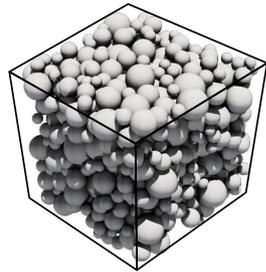
micro-

Macro-scale

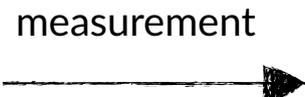
auxiliary problem

thermodynamics-based neural nets

BVP

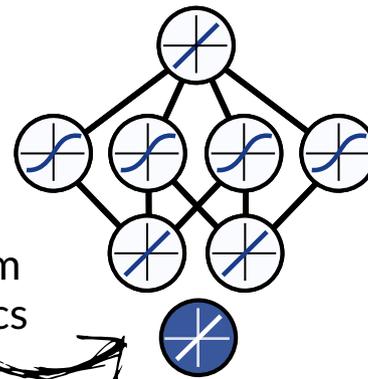


unit cell



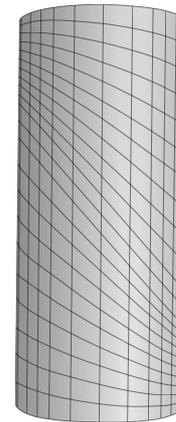
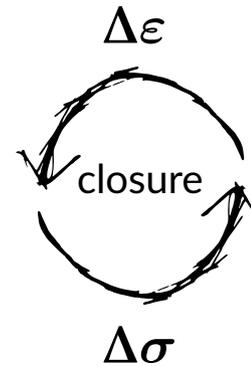
non-equilibrium
thermodynamics

Coleman & Gurtin, 1968



Masi et al. 2021,
Raissi et al, 2021,
Chinesta et al. 2021

**massive speeds-up,
consistent, good generalisation**



$$\begin{cases} \nabla \sigma + \rho \mathbf{b} = \rho \ddot{\mathbf{u}} \\ \epsilon = \text{Sym}(\nabla \mathbf{u}) \end{cases}$$

I) Background setting

- a) Non-equilibrium thermodynamics

II) Thermodynamics-based artificial neural networks

- a) Methods
- b) Discovery of internal state variables
- c) Experiments

III) Multiscale computing

IV) Neural equations from small data

- a) Data scarcity and sparsity
- b) Methods
- c) Experiments

I) Background setting

- a) Non-equilibrium thermodynamics

II) Thermodynamics-based artificial neural networks

- a) Methods
- b) Discovery of internal state variables
- c) Experiments

III) Multiscale computing

IV) Neural equations from small data

- a) Data scarcity and sparsity
- b) Methods
- c) Experiments

Non-equilibrium thermodynamics

Coleman and Gurtin, 1967

- Balance of energy and entropy (local)

$$\dot{e} = p^i - \nabla \cdot \mathbf{q} + r, \quad (1)$$

$$\gamma = \dot{\eta} + \nabla \cdot \left(\frac{\mathbf{q}}{T} \right) - \frac{r}{T} \geq 0 \quad (2)$$

internal energy density e , heat flux, source \mathbf{q}, r , power of internal forces $p^i = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}}$
entropy density $\dot{\eta}$, absolute temperature T , dissipation rate γ

- Balance of energy and entropy (local)

$$\dot{e} = p^i - \nabla \cdot \mathbf{q} + r, \quad (1) \quad \gamma = \dot{\eta} + \nabla \cdot \left(\frac{\mathbf{q}}{T} \right) - \frac{r}{T} \geq 0 \quad (2)$$

internal energy density e , heat flux, source \mathbf{q}, r , power of internal forces $p^i = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}}$
entropy density $\dot{\eta}$, absolute temperature T , dissipation rate γ

- Clausius-Duhem inequality

$$\gamma = p^i - \dot{e} + T\dot{\eta} - \mathbf{q} \cdot \frac{\nabla T}{T} = p^i - (\dot{\psi} + \dot{T}\eta) - \mathbf{q} \cdot \frac{\nabla T}{T} \geq 0 \quad (3)$$

$$d = p^i - (\dot{\psi} + \dot{T}\eta) \geq 0 \quad (4)$$

Helmholtz free energy $\psi = e - T\eta$, dissipation rate d

- Balance of energy and entropy (local)

$$\dot{e} = p^i - \nabla \cdot \mathbf{q} + r, \quad (1) \quad \gamma = \dot{\eta} + \nabla \cdot \left(\frac{\mathbf{q}}{T} \right) - \frac{r}{T} \geq 0 \quad (2)$$

internal energy density e , heat flux, source \mathbf{q}, r , power of internal forces $p^i = \boldsymbol{\sigma} : \dot{\boldsymbol{\epsilon}}$
 entropy density $\dot{\eta}$, absolute temperature T , dissipation rate γ

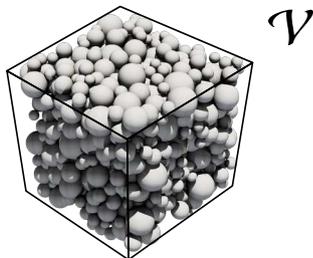
- Clausius-Duhem inequality

$$\gamma = p^i - \dot{e} + T\dot{\eta} - \mathbf{q} \cdot \frac{\nabla T}{T} = p^i - (\dot{\psi} + \dot{T}\eta) - \mathbf{q} \cdot \frac{\nabla T}{T} \geq 0 \quad (3)$$

$$d = p^i - (\dot{\psi} + \dot{T}\eta) \geq 0 \quad (4)$$

Helmholtz free energy $\psi = e - T\eta$, dissipation rate d

- Volume average Clausius-Duhem inequality



$$\langle d \rangle = \langle p^i \rangle - \langle \dot{\psi} \rangle - \langle \dot{T}\eta \rangle \geq 0 \quad (5)$$

$$\langle p^i \rangle = \langle \boldsymbol{\sigma} \rangle : \langle \dot{\boldsymbol{\epsilon}} \rangle$$

$$\langle \psi \rangle = \frac{1}{|\mathcal{V}|} \int_{\mathcal{V}} \psi \, dy$$

(periodic displ. & anti-periodic tract.)

Non-equilibrium thermodynamics

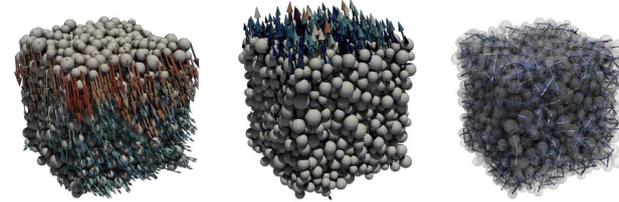
Coleman and Gurtin, 1967

4

- State functions and variables

$$\psi = \hat{\psi}(\mathbb{X}) \quad \{\varepsilon, T, \mathbf{z}, \rho, \cdot\} \subset \mathbb{X} \quad (6)$$

internal state variables $\mathbf{z} = \{z_1, z_2, \dots, z_p\}$



Non-equilibrium thermodynamics

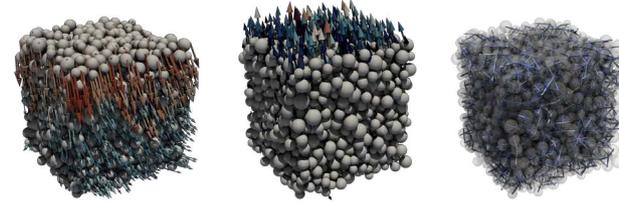
Coleman and Gurtin, 1967

4

- State functions and variables

$$\psi = \hat{\psi}(\mathbb{X}) \quad \{\varepsilon, T, \mathbf{z}, \rho, \cdot\} \subset \mathbb{X} \quad (6)$$

internal state variables $\mathbf{z} = \{z_1, z_2, \dots, z_p\}$



- Evolution equations for the internal variables (general)

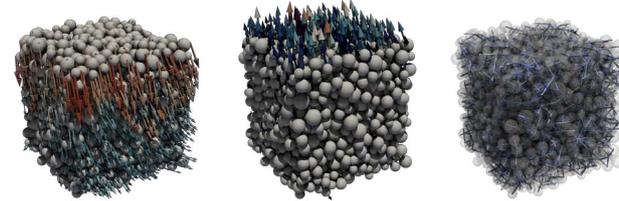
$$\dot{\mathbf{z}} = \mathbf{f}(\mathbb{X}) \quad (7)$$

(special cases: Onsager-Casimir relationships, Generalised Standard Materials, etc.)

Onsager 1931, Moreau 1973, Halphen, Son Nguyen, 1975, de Saxcé 1998

- State functions and variables $\psi = \hat{\psi}(\mathbb{X}) \quad \{\varepsilon, T, \mathbf{z}, \rho, \cdot\} \subset \mathbb{X} \quad (6)$

internal state variables $\mathbf{z} = \{z_1, z_2, \dots, z_p\}$



- Evolution equations for the internal variables (general) $\dot{\mathbf{z}} = \mathbf{f}(\mathbb{X}) \quad (7)$

(special cases: Onsager-Casimir relationships, Generalised Standard Materials, etc.)

Onsager 1931, Moreau 1973, Halphen, Son Nguyen, 1975, de Saxcé 1998

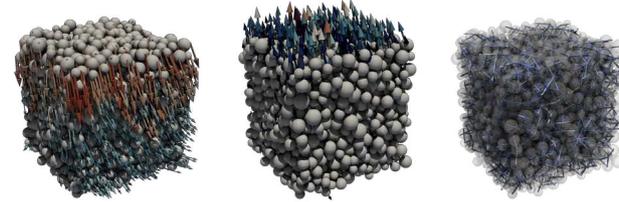
- Substituting the time derivative of $\hat{\psi}$ into the CD inequality yields

$$\left(\boldsymbol{\sigma} - \frac{\partial \hat{\psi}}{\partial \boldsymbol{\varepsilon}} \right) : \dot{\boldsymbol{\varepsilon}} + \left(\frac{\partial \hat{\psi}}{\partial T} + \eta \right) \dot{T} + \frac{\partial \hat{\psi}}{\partial \mathbf{z}} \cdot \dot{\mathbf{z}} - d = 0 \quad \forall \dot{\boldsymbol{\varepsilon}}, \dot{T}, \dot{\mathbf{z}} \quad (8)$$

$$d \geq 0$$

- State functions and variables $\psi = \hat{\psi}(\mathbb{X}) \quad \{\varepsilon, T, \mathbf{z}, \rho, \cdot\} \subset \mathbb{X} \quad (6)$

internal state variables $\mathbf{z} = \{\mathbf{z}_1, \mathbf{z}_2, \dots, \mathbf{z}_p\}$



- Evolution equations for the internal variables (general) $\dot{\mathbf{z}} = \mathbf{f}(\mathbb{X}) \quad (7)$

(special cases: Onsager-Casimir relationships, Generalised Standard Materials, etc.)

Onsager 1931, Moreau 1973, Halphen, Son Nguyen, 1975, de Saxcé 1998

- Substituting the time derivative of $\hat{\psi}$ into the CD inequality yields

$$\left(\boldsymbol{\sigma} - \frac{\partial \hat{\psi}}{\partial \boldsymbol{\varepsilon}} \right) : \dot{\boldsymbol{\varepsilon}} + \left(\frac{\partial \hat{\psi}}{\partial T} + \eta \right) \dot{T} + \frac{\partial \hat{\psi}}{\partial \mathbf{z}} \cdot \dot{\mathbf{z}} - d = 0 \quad \forall \dot{\boldsymbol{\varepsilon}}, \dot{T}, \dot{\mathbf{z}} \quad (8)$$

$$d \geq 0$$

- Thermodynamics-admissible material processes

$$\boldsymbol{\sigma} = \frac{\partial \hat{\psi}}{\partial \boldsymbol{\varepsilon}}(\mathbb{X}), \quad \eta = -\frac{\partial \hat{\psi}}{\partial T}(\mathbb{X}) \quad d = \mathbb{Z} \cdot \dot{\mathbf{z}} \geq 0 \quad (9)$$

$$\mathbb{Z} \equiv -\partial_{\mathbf{z}} \hat{\psi}(\mathbb{X})$$

I) Background setting

- a) Non-equilibrium thermodynamics

II) Thermodynamics-based artificial neural networks

- a) Methods
- b) Discovery of internal state variables
- c) Experiments

III) Multiscale computing

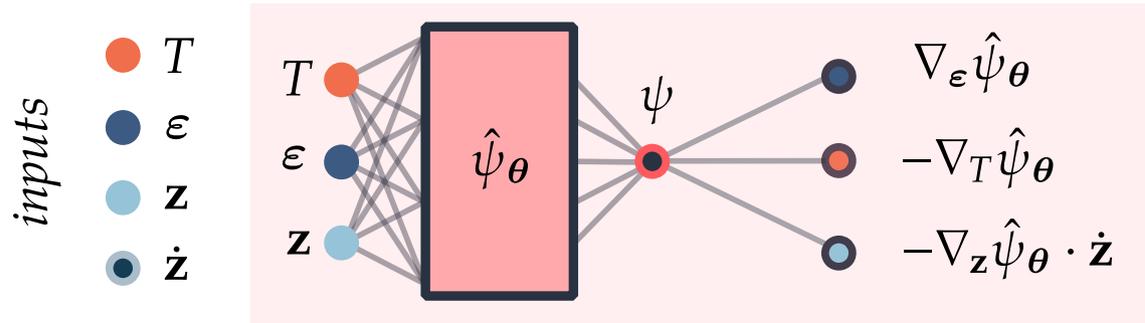
IV) Neural equations from small data

- a) Data scarcity and sparsity
- b) Methods
- c) Experiments

Thermodynamics-based Neural Nets

Masi et al. JMPS, 2021; Masi & Stefanou. JMPS, 2023

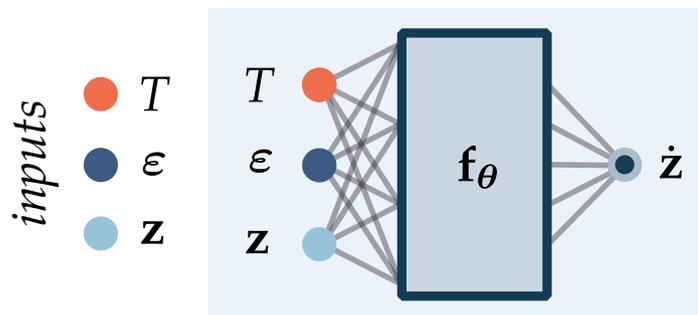
I) Free-energy density network $\hat{\psi}_\theta$



$$\theta^* = \arg \min_{\theta} \left(\underbrace{\ell(\psi, \hat{\psi}_\theta)}_{\text{loss in } \psi} + \underbrace{\ell(\sigma, \nabla_\epsilon \hat{\psi}_\theta) + \ell(\eta, -\nabla_T \hat{\psi}_\theta) + \ell(d, -\nabla_z \hat{\psi}_\theta \cdot \dot{\mathbf{z}})}_{\text{loss in } \nabla \psi} + \underbrace{\langle -\nabla_z \hat{\psi}_\theta \cdot \dot{\mathbf{z}} \rangle}_{\text{reg.}} \right)$$

$$\langle a \rangle = \max(0, a)$$

II) Evolution equation network \mathbf{f}_θ

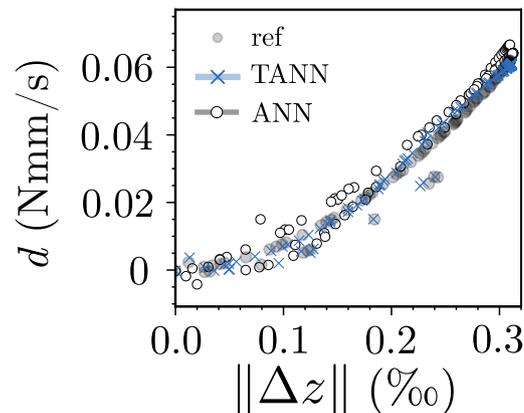
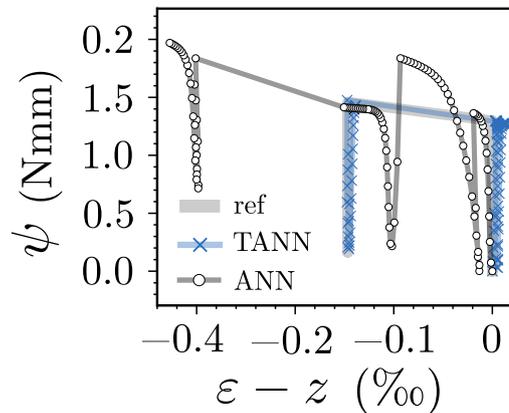
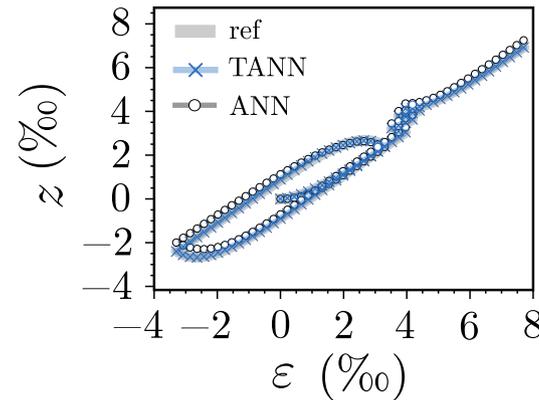
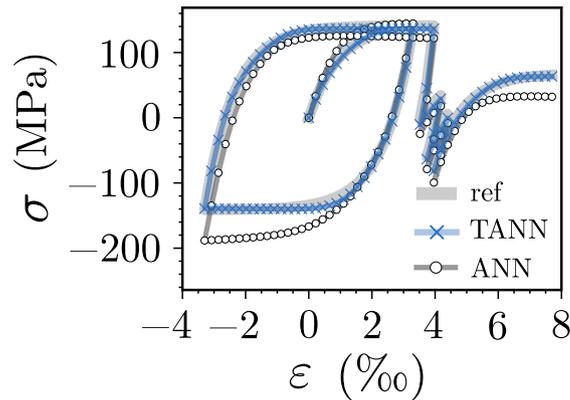
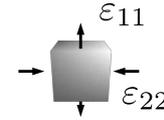


$$\theta^* = \arg \min_{\theta} \left(\underbrace{\ell(\dot{\mathbf{z}}, \mathbf{f}_\theta)}_{\text{loss in } \dot{\mathbf{z}}} \right)$$

Data-driven versus physics-driven

ANN versus TANN: why thermodynamics is important!

▷ Biaxial strain-driven ratcheting



▷ Data-driven (ANN)

- Accuracy
- Energy is not conserved
- Dissipation rate is negative

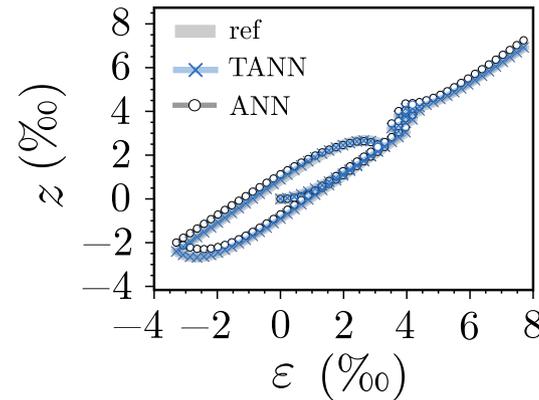
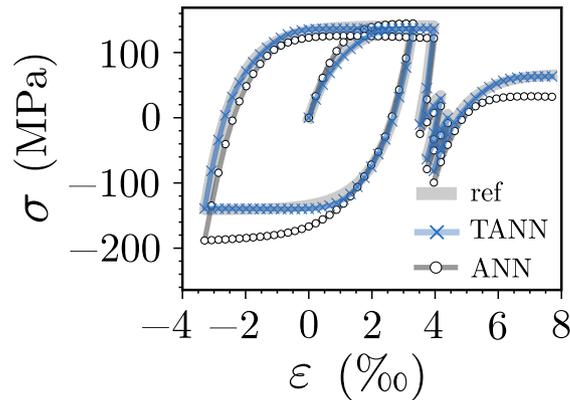
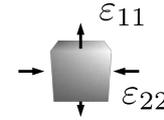
▷ Physics-driven (TANN)

- + Accuracy
- + Energy is conserved
- + Dissipation is positive-semi def.

Data-driven versus physics-driven

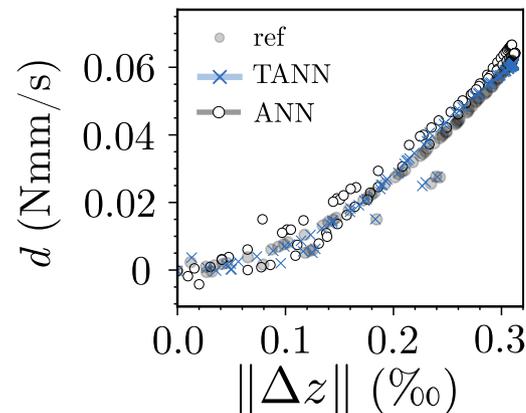
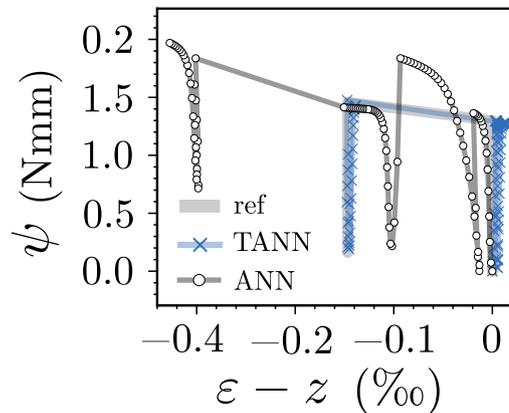
ANN versus TANN: why thermodynamics is important!

▷ Biaxial strain-driven ratcheting



▷ Data-driven (ANN)

- Accuracy
- Energy is not conserved
- Dissipation rate is negative



▷ Physics-driven (TANN)

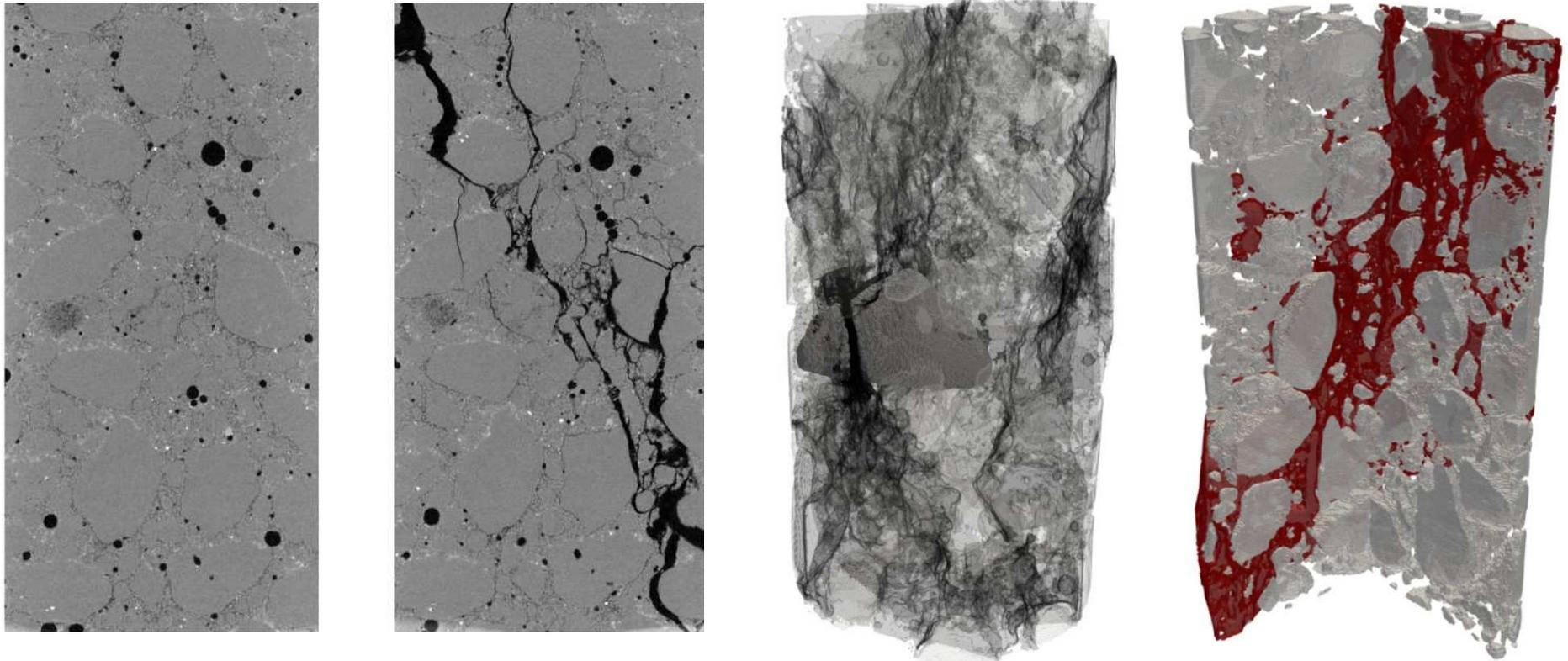
- + Accuracy
- + Energy is conserved
- + Dissipation is positive-semi def.

▷ High performances for a wide spectrum of rate-(in)dependent inelastic behaviours

Internal state variables and where to find them

Internal state variables and where to find them

Laboratory experiment: concrete under triaxial compression (x-ray tomography)



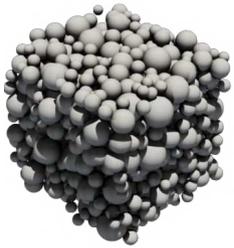
Discovery of internal state variables

Masi & Stefanou. CMAME, 2022

9

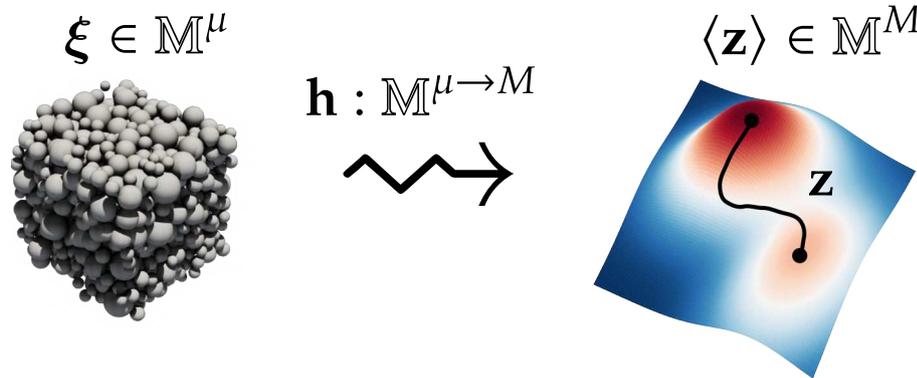
- Internal coordinates micro-related quantities describing internal mechanisms (observable – degrees of freedom)

$$\xi \in \mathbb{M}^\mu$$



Discovery of internal state variables

- Internal coordinates micro-related quantities describing internal mechanisms (observable – degrees of freedom)

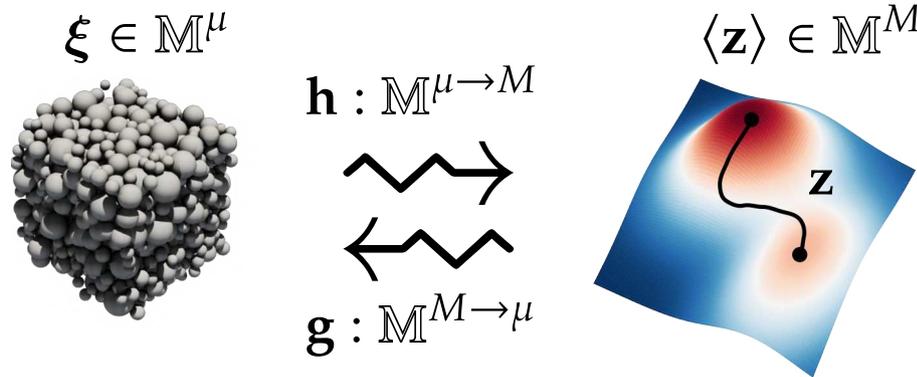


- Discovered internal variables: latent variables of the internal coordinates

$$\mathbf{z} = \mathbf{h}(\xi) \quad \text{such that} \quad \psi = \hat{\psi}(\mathbb{X})$$

Discovery of internal state variables

- Internal coordinates micro-related quantities describing internal mechanisms (observable – degrees of freedom)



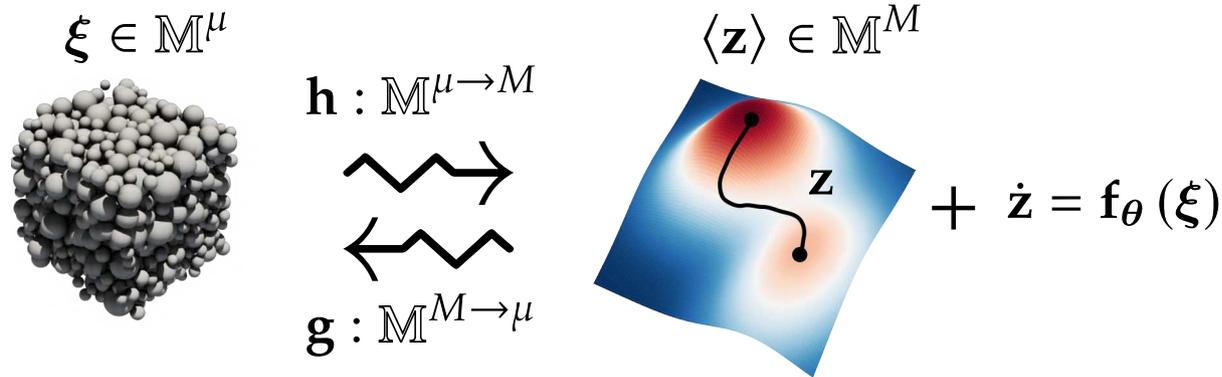
- Discovered internal variables: latent variables of the internal coordinates

$$\mathbf{z} = \mathbf{h}(\xi) \quad \text{such that} \quad \psi = \hat{\psi}(\mathbb{X})$$

$$\tilde{\xi} = \mathbf{g}(\mathbf{z}) \quad \text{pseudo-inverse transform}$$

Discovery of internal state variables

- Internal coordinates micro-related quantities describing internal mechanisms (observable – degrees of freedom)



- Discovered internal variables: latent variables of the internal coordinates

$$\mathbf{z} = \mathbf{h}(\xi) \quad \text{such that} \quad \psi = \hat{\psi}(\mathbb{X})$$

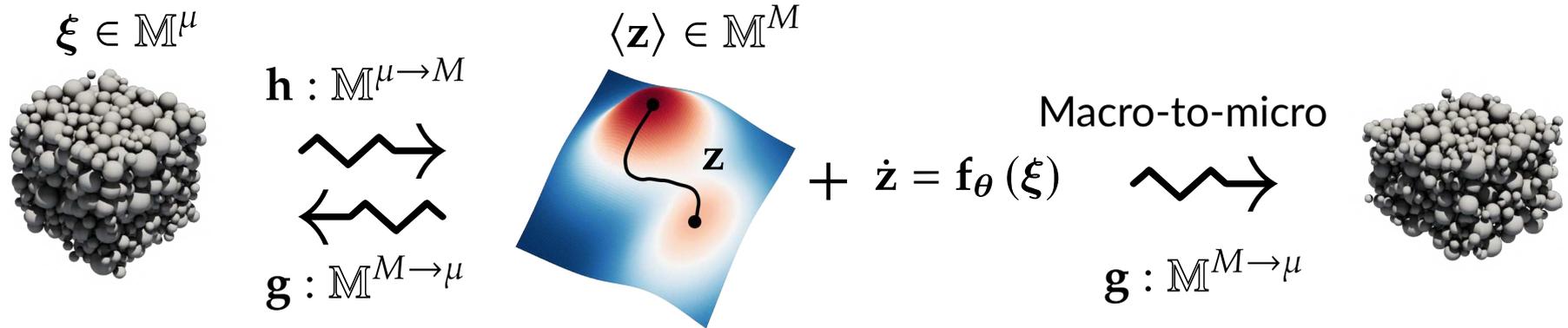
$$\tilde{\xi} = \mathbf{g}(\mathbf{z}) \quad \text{pseudo-inverse transform}$$

- Discovered evolution equations: dynamics of the internal coordinates

$$\dot{\mathbf{z}} = \mathbf{f}_\theta(\mathbb{X}) \quad \iff \quad \dot{\mathbf{z}} = \frac{d\mathbf{h}}{dt}(\xi) = \nabla_\xi \mathbf{h}(\xi) \cdot \dot{\xi}$$

Discovery of internal state variables

- Internal coordinates micro-related quantities describing internal mechanisms (observable – degrees of freedom)



- Discovered internal variables: latent variables of the internal coordinates

$$\mathbf{z} = \mathbf{h}(\xi) \quad \text{such that} \quad \psi = \hat{\psi}(\mathbb{X})$$

$$\tilde{\xi} = \mathbf{g}(\mathbf{z}) \quad \text{pseudo-inverse transform}$$

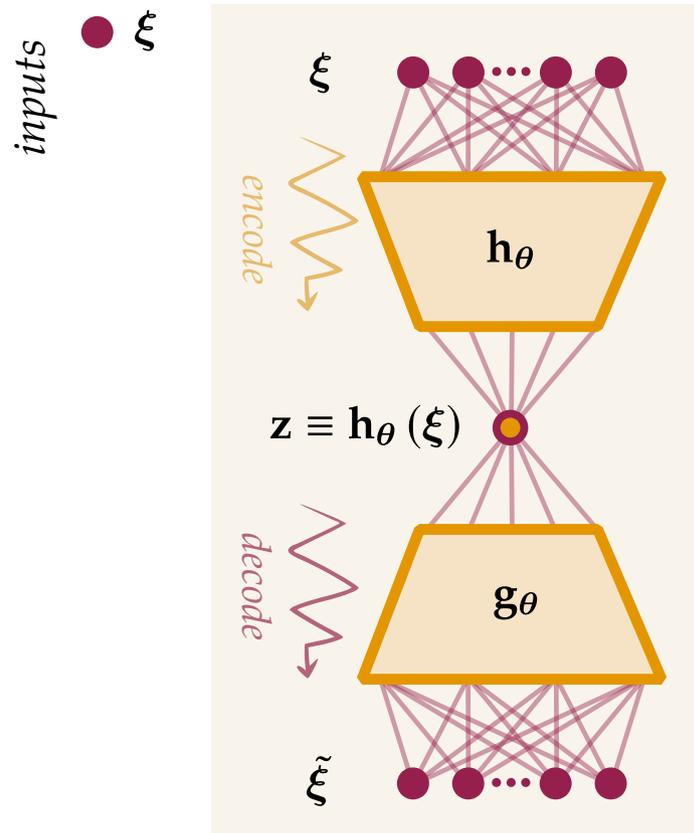
- Discovered evolution equations: dynamics of the internal coordinates

$$\dot{\mathbf{z}} = \mathbf{f}_\theta(\mathbb{X}) \quad \iff \quad \dot{\mathbf{z}} = \frac{d\mathbf{h}}{dt}(\xi) = \nabla_\xi \mathbf{h}(\xi) \cdot \dot{\xi}$$

$$\dot{\xi} = \frac{d\mathbf{g}}{dt}(\mathbf{z}) = \nabla_{\mathbf{z}} \mathbf{g}(\mathbf{z}) \cdot \dot{\mathbf{z}}$$

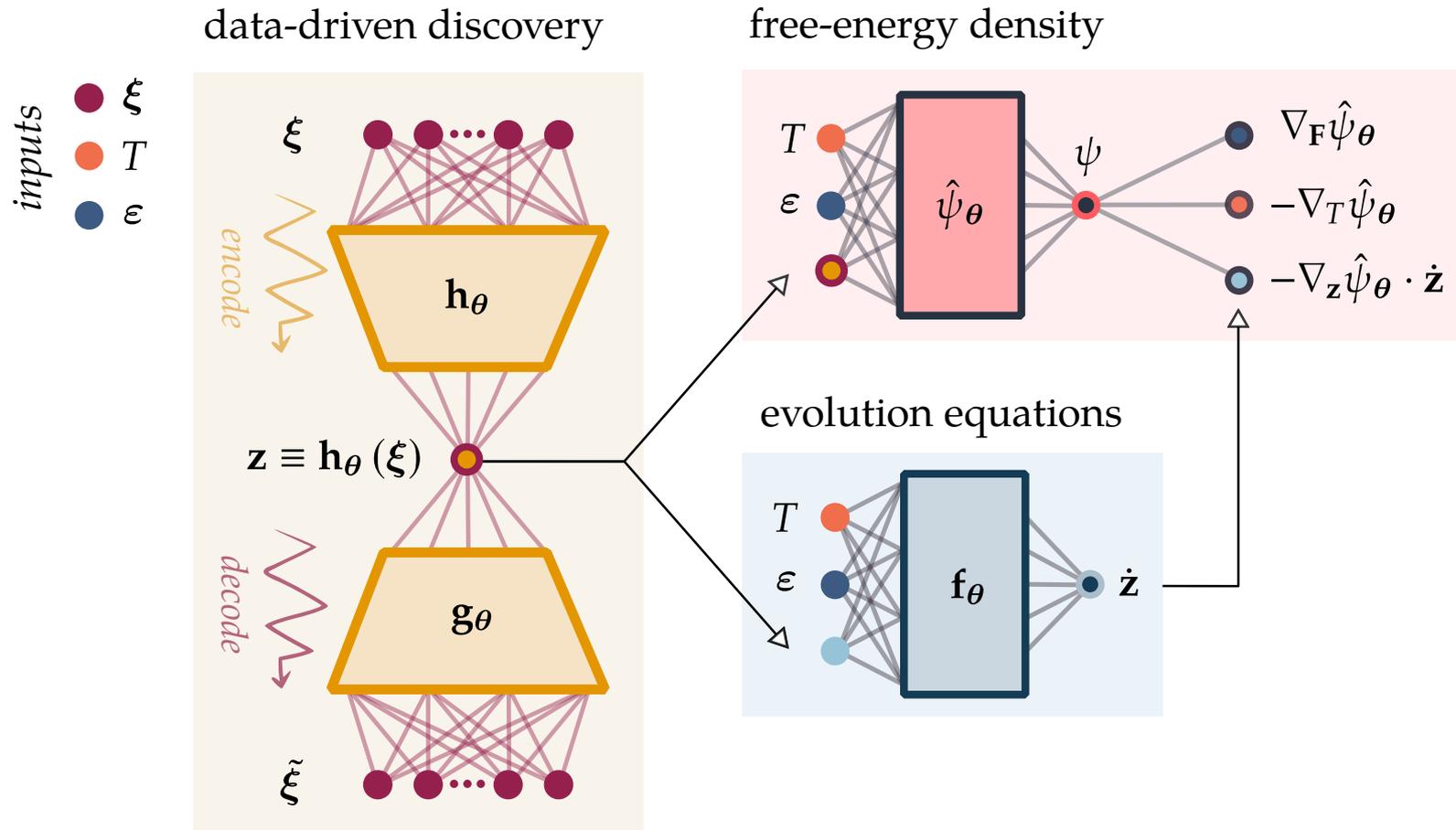
Discovery of internal state variables

data-driven discovery



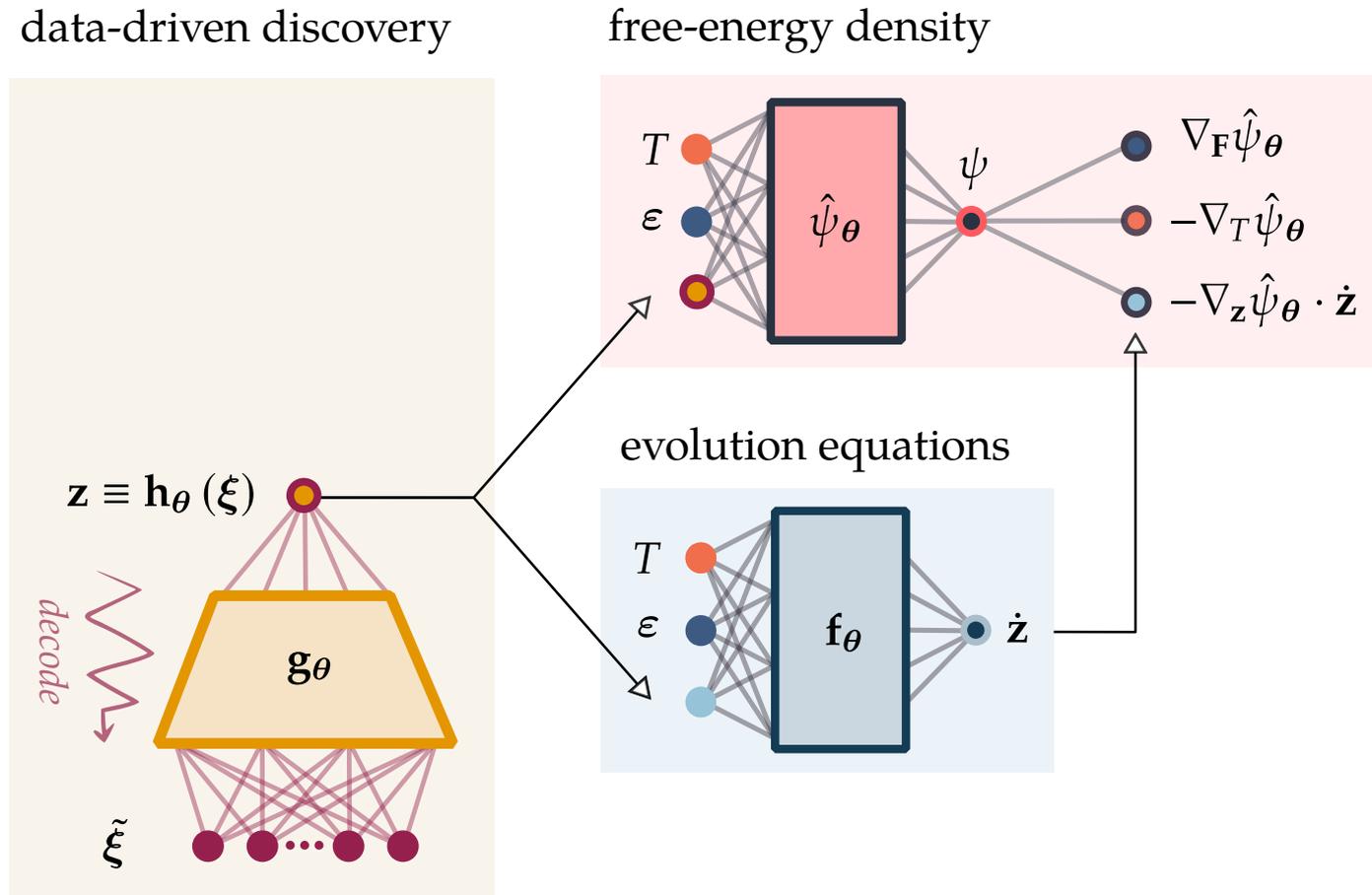
$$\theta^* = \arg \min_{\theta} \left(\underbrace{\ell(\xi, g_\theta(h_\theta(\xi)))}_{\text{reconstruction loss}} \right)$$

Discovery of internal state variables



$$\theta^* = \arg \min_{\theta} \left(\underbrace{\ell(\xi, g_\theta(h_\theta(\xi)))}_{\text{reconstruction loss}} + \underbrace{\ell(\dot{\xi}, \nabla_z g_\theta \cdot \dot{z})}_{\text{loss in } \dot{\xi}} + \underbrace{\ell(\psi, \hat{\psi}_\theta)}_{\text{loss in } \psi} \right. \\ \left. + \underbrace{\ell(\mathbf{P}, \nabla_F \hat{\psi}_\theta) + \ell(\eta, -\nabla_T \hat{\psi}_\theta) + \ell(d, -\nabla_z \hat{\psi}_\theta \cdot \dot{z})}_{\text{loss in } \nabla \psi} + \underbrace{\langle -\nabla_z \hat{\psi}_\theta \cdot \dot{z} \rangle}_{\text{reg.}} \right)$$

Discovery of internal state variables



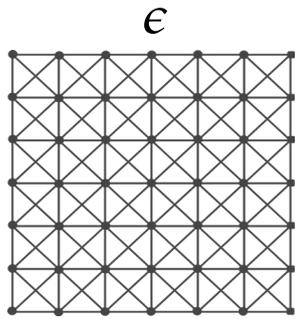
Constitutive equations: $\dot{\mathbf{z}}(t), \boldsymbol{\sigma}(t) = \mathcal{T}_{\mathcal{ANN}}(T(t), \boldsymbol{\epsilon}(t), \mathbf{z}(t)) \quad \forall t$

Initial Value Problem: $\mathbf{z}(t) = \int_{t_0}^t \mathbf{f}_\theta(T(t), \boldsymbol{\epsilon}(t), \mathbf{z}(t)) + \mathbf{z}(t_0)$
 $\mathbf{z}(t_0) = \mathbf{h}_\theta(\boldsymbol{\xi}(t_0))$

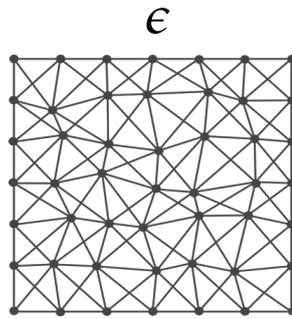
Downscaling - "localisation": $\boldsymbol{\xi}(t) = \mathbf{g}_\theta(\mathbf{z}(t))$

Experiments: lattice materials

Heterogeneous lattice structures (auxiliary problem)



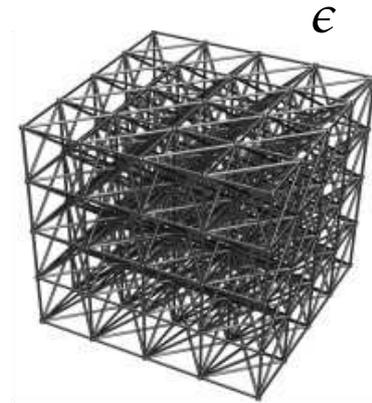
2D regular



2D irregular



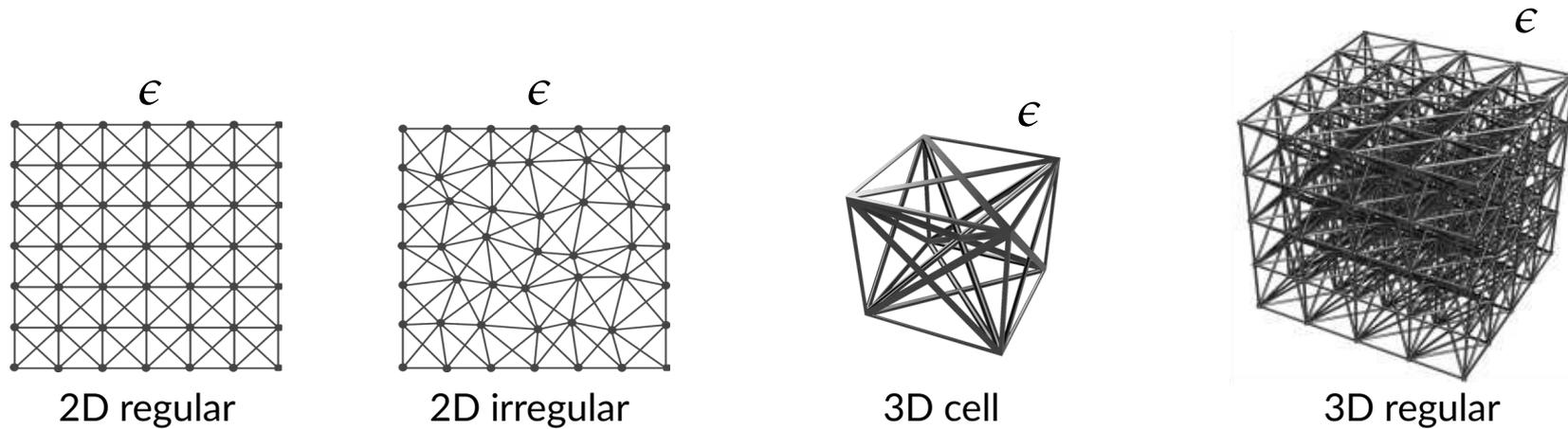
3D cell



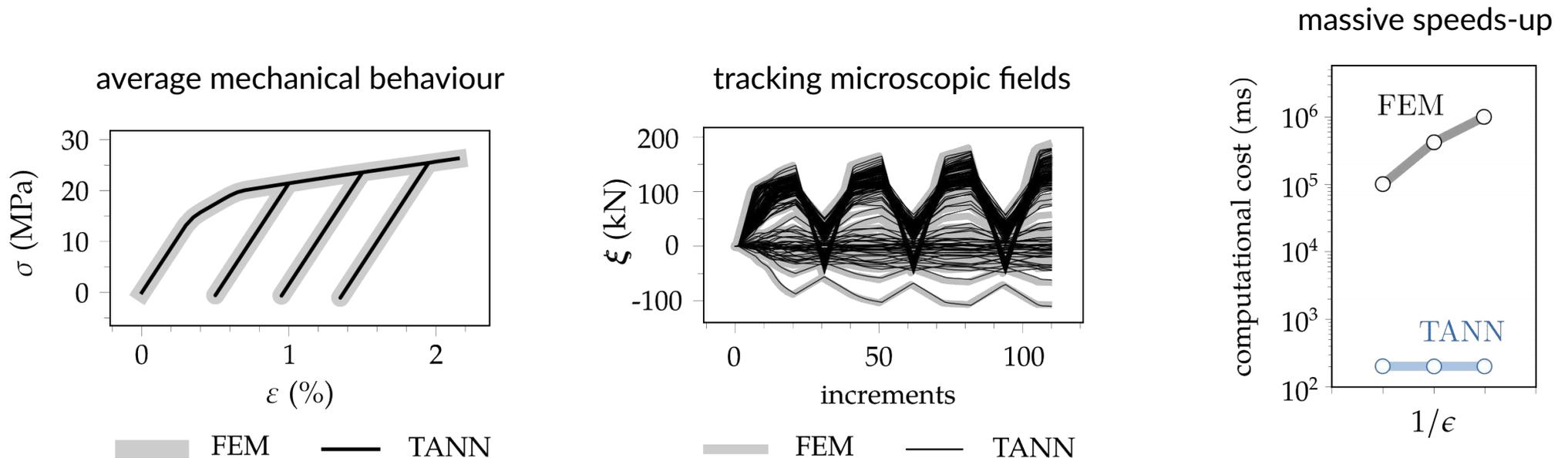
3D regular

Experiments: lattice materials

Heterogeneous lattice structures (auxiliary problem)



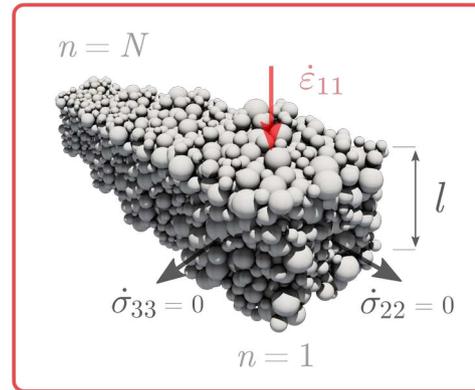
Replacing the solution of the auxiliary problem (micromechanical FEM versus TANN)



Experiments: granular matter

Experiments: granular matter

Stochastic representation (random media) – statistical ensemble approach



statistical
elementary volumes
(SEV)

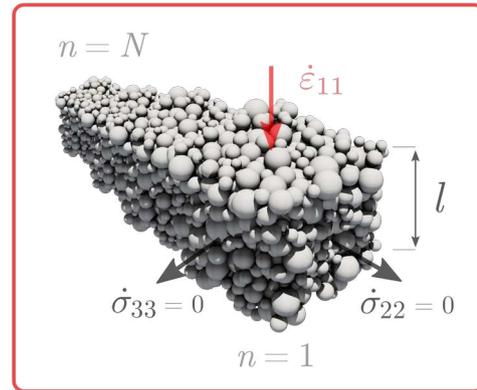


representative
elementary volume
(REV)

Ostoja-Starzewski 2006
Papachristos et al. 2023
Nguyen, 2021

Experiments: granular matter

Stochastic representation (random media) – statistical ensemble approach

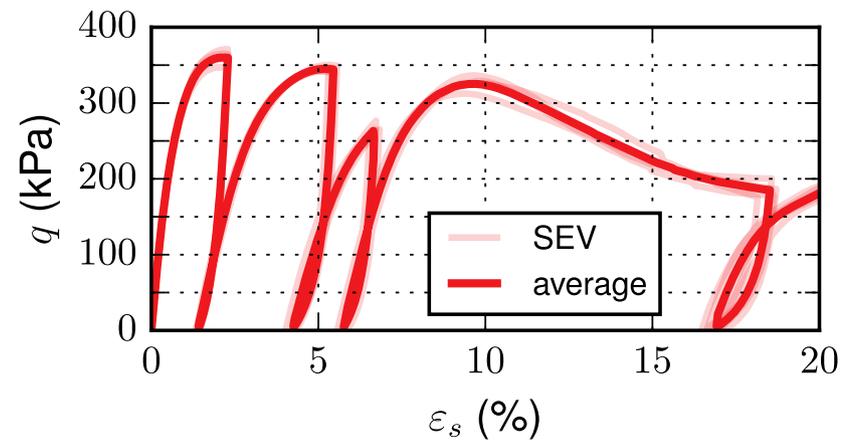
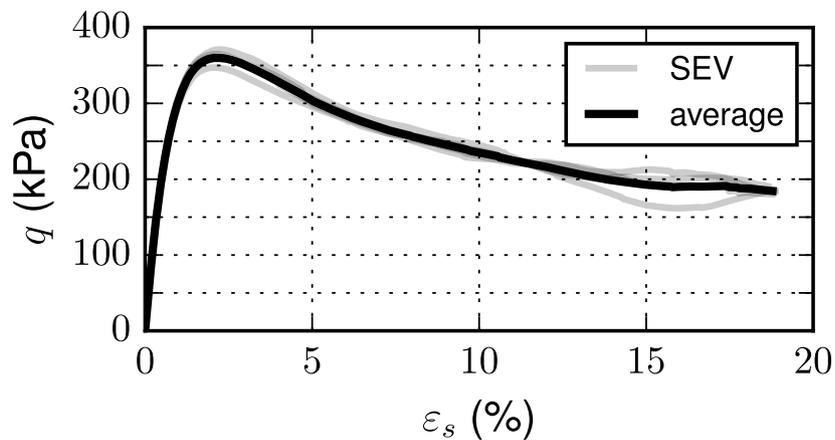
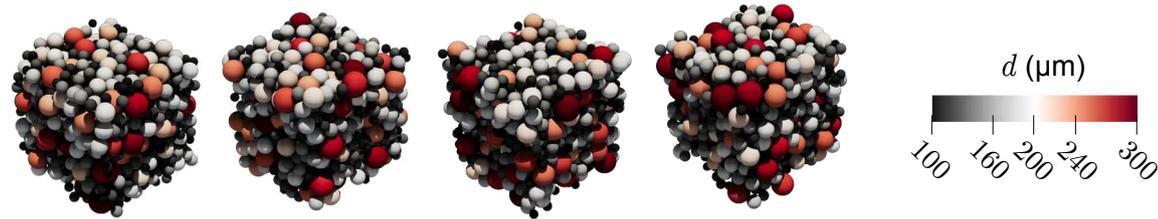


statistical elementary volumes (SEV)



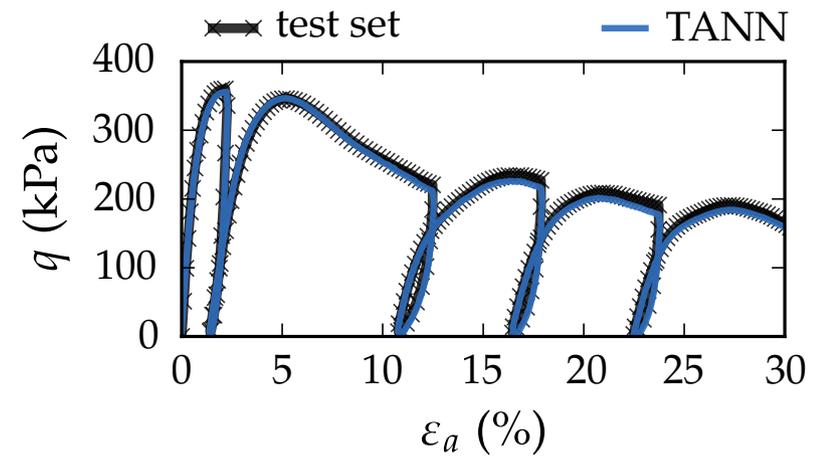
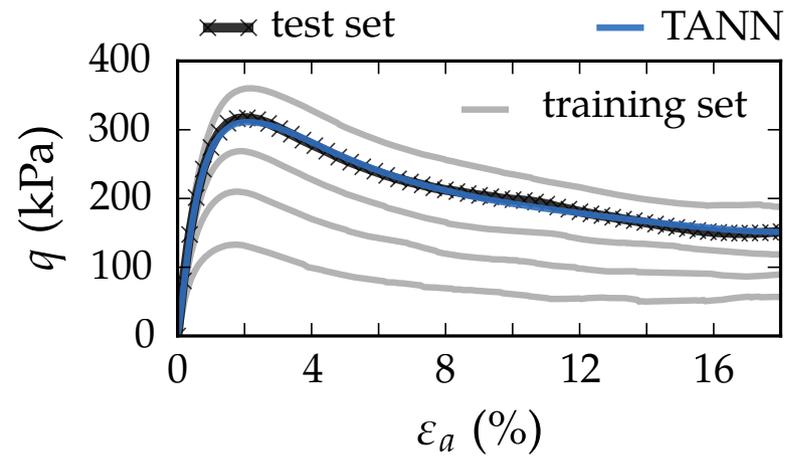
representative elementary volume (REV)

Ostoja-Starzewski 2006
Papachristos et al. 2023
Nguyen, 2021



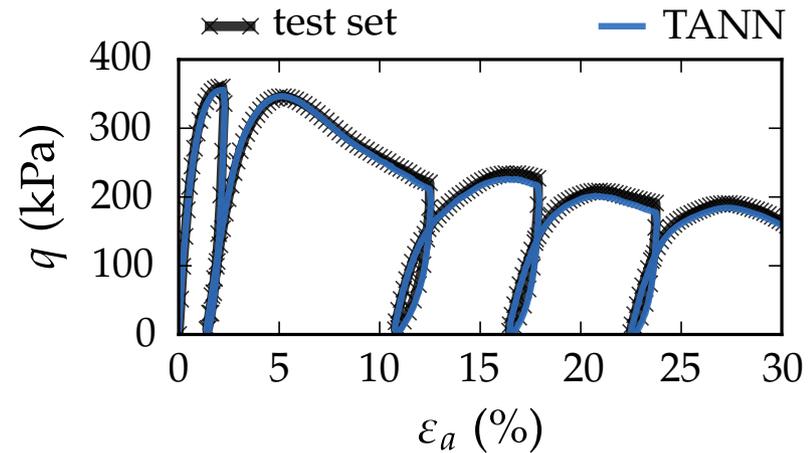
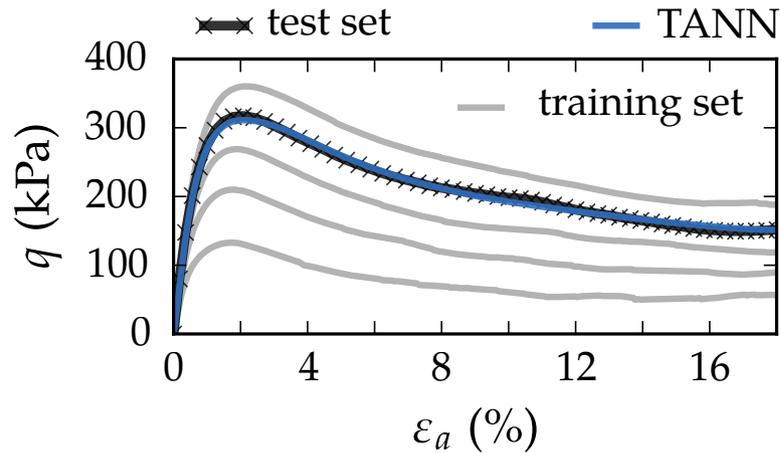
Experiments: granular matter

Predicting unobserved behaviours

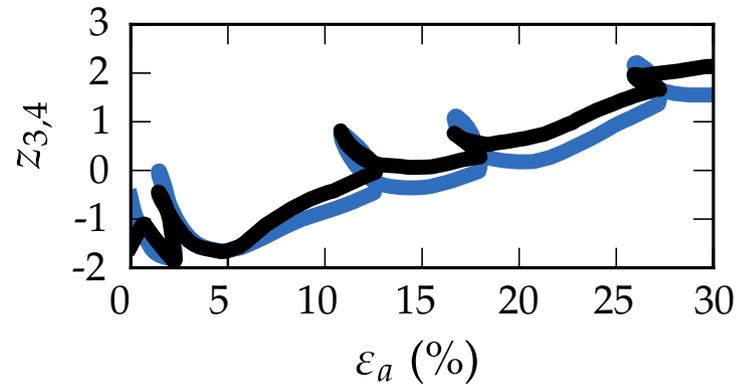
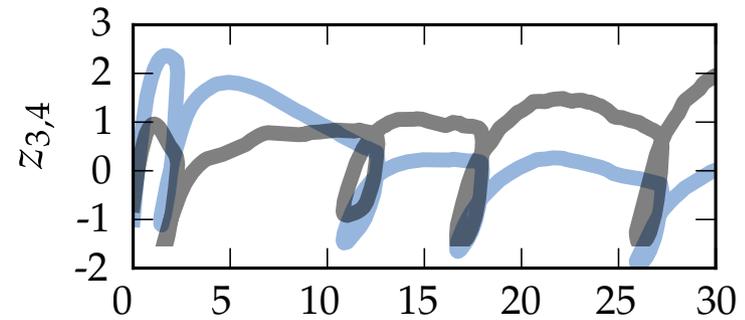
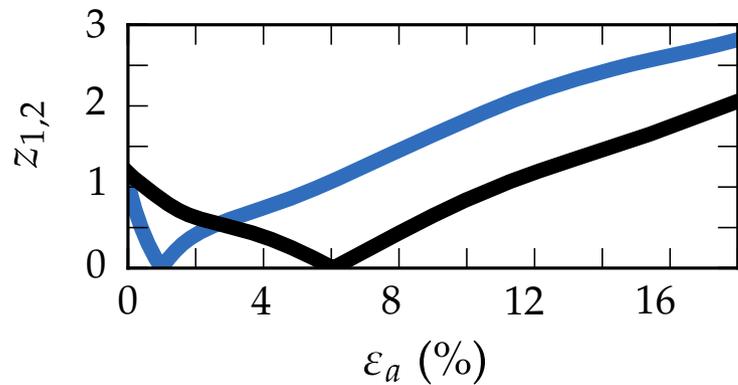
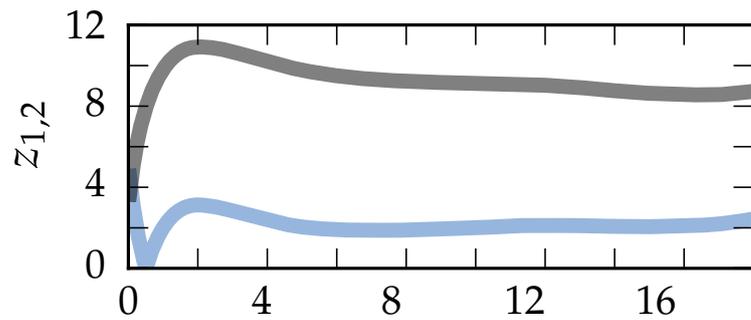


Experiments: granular matter

Predicting unobserved behaviours



- micro-to-Macro: 4 internal state variables \sim topological descriptors



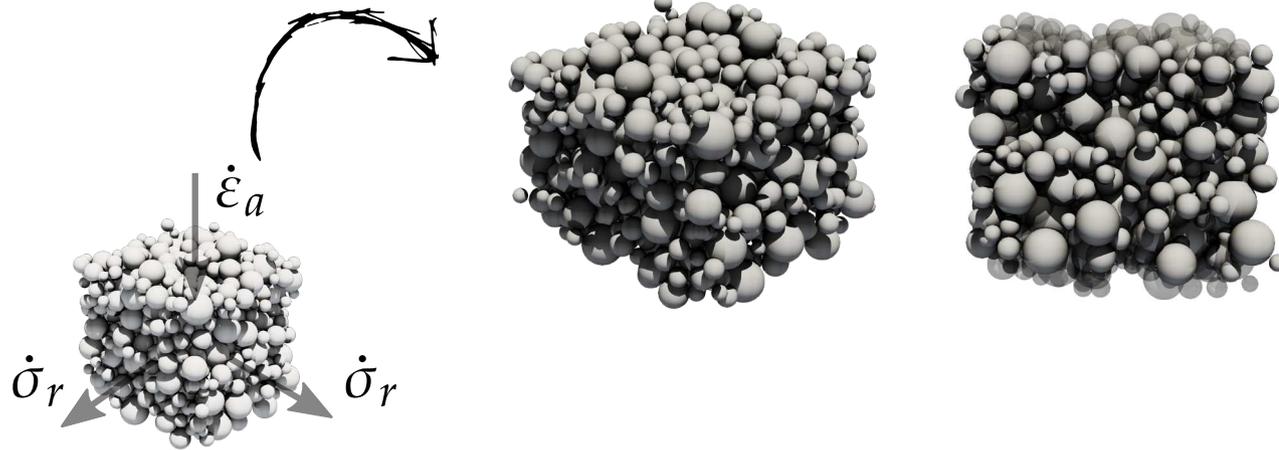
On-the-fly microstructure tracking

- Reconstructing micromechanical fields (data-driven localisation)

On-the-fly microstructure tracking

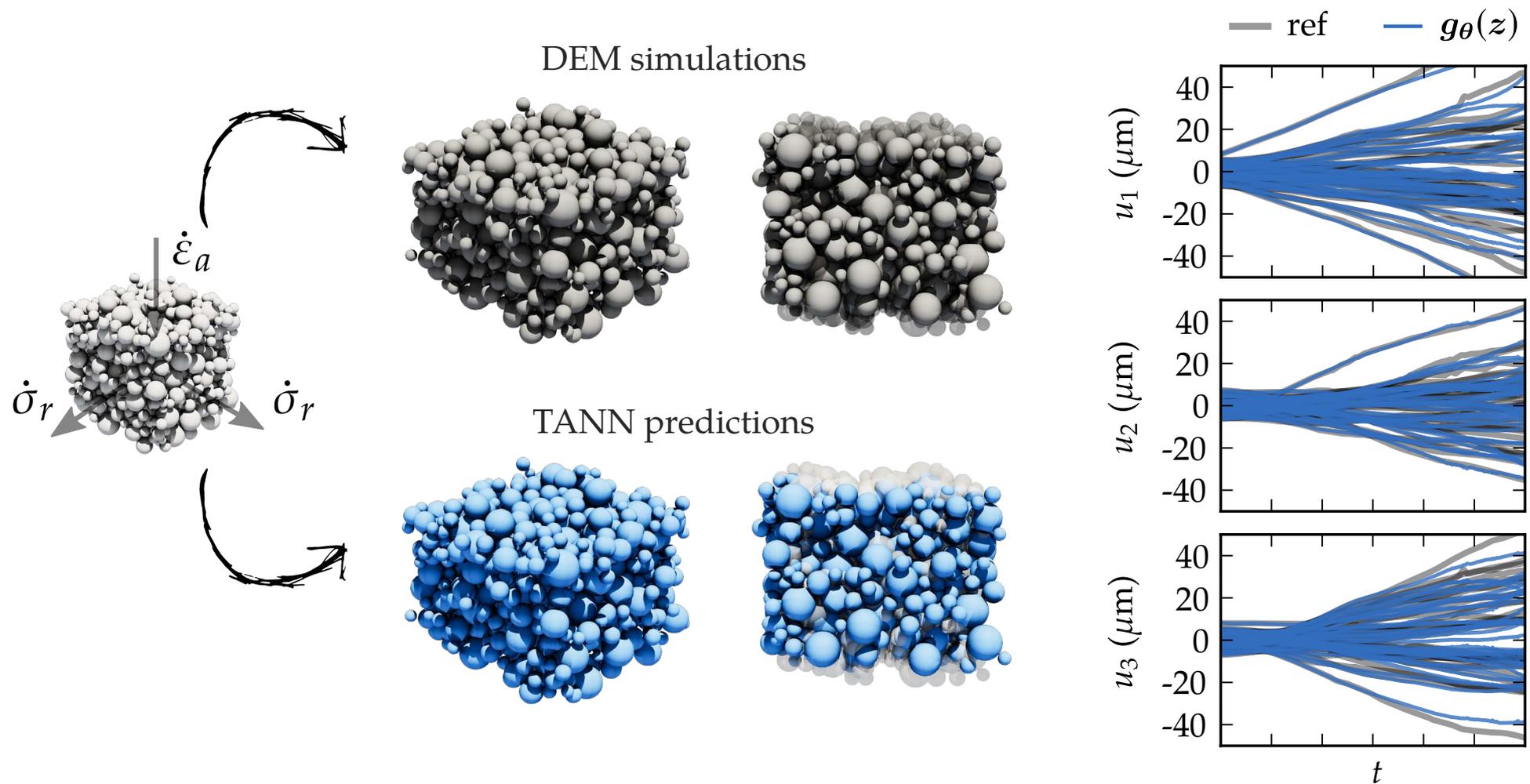
- Reconstructing micromechanical fields (data-driven localisation)

DEM simulations



On-the-fly microstructure tracking

- Reconstructing micromechanical fields (data-driven localisation)



▷ Massive computational accelerations ($\times 10^3$) – DEM/TANN: 30 minutes / $< 0.3\text{s}$

I) Background setting

- a) Non-equilibrium thermodynamics

II) Thermodynamics-based artificial neural networks

- a) Methods
- b) Discovery of internal state variables
- c) Experiments

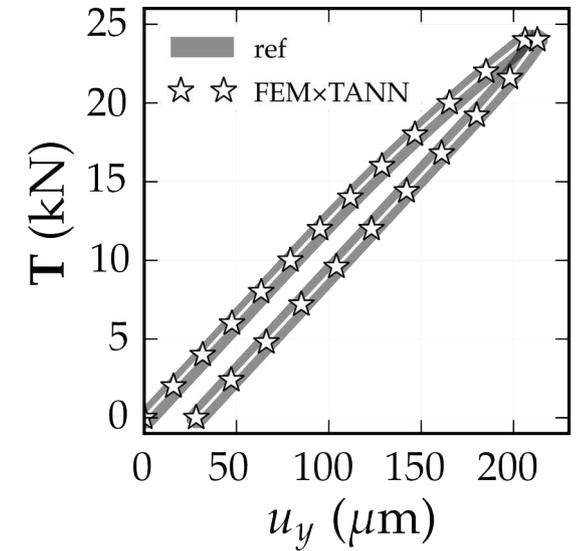
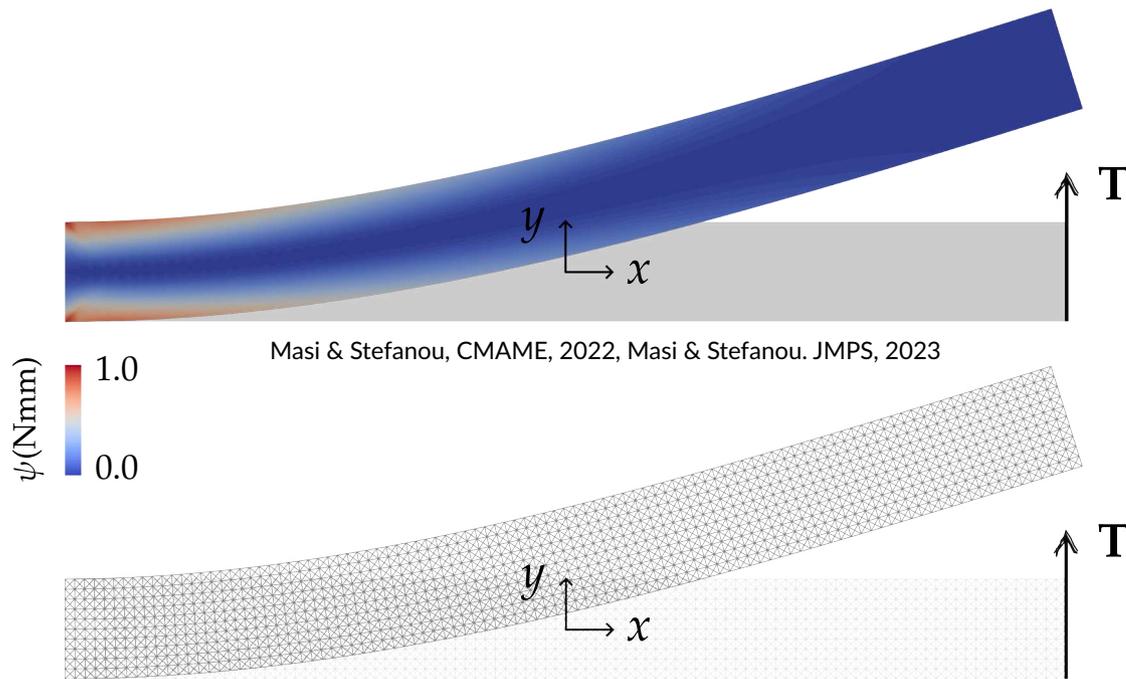
III) Multiscale computing

IV) Neural equations from small data

- a) Data scarcity and sparsity
- b) Methods
- c) Experiments

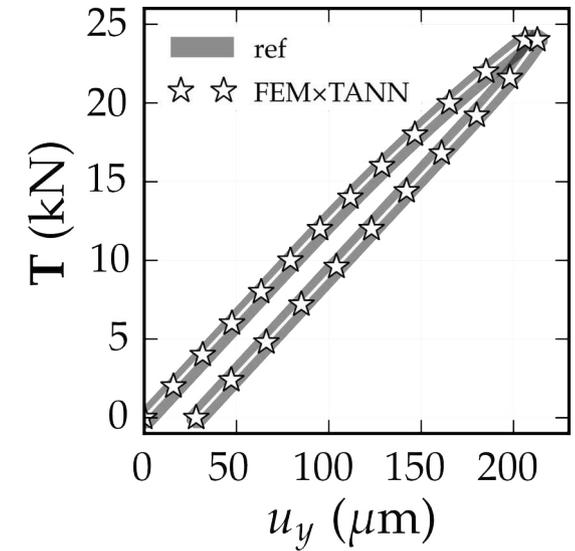
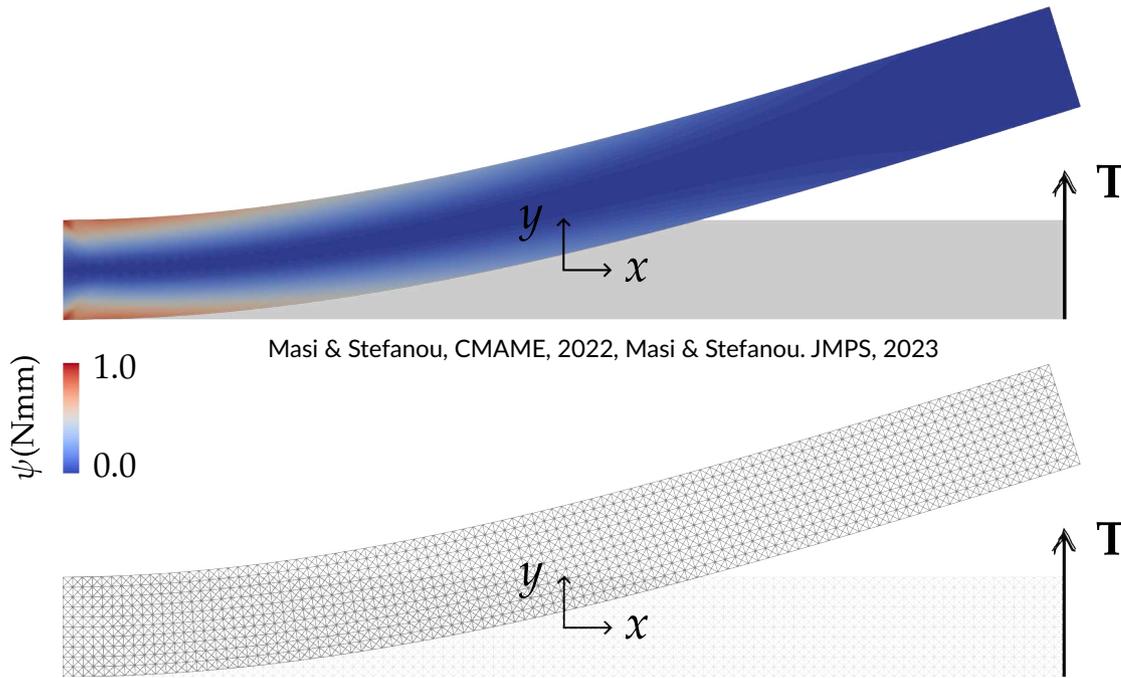
FEM×TANN and numerical homogenisation

Double-scale homogenisation scheme: rigorous (asymptotic) and fast

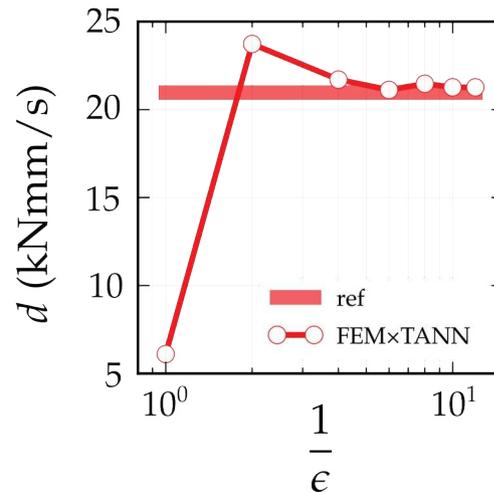
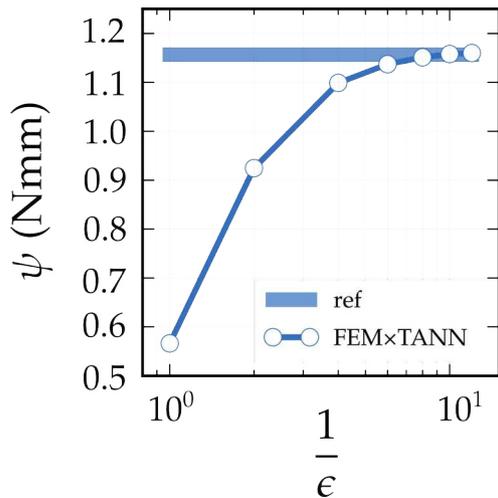


FEM×TANN and numerical homogenisation

Double-scale homogenisation scheme: rigorous (asymptotic) and fast

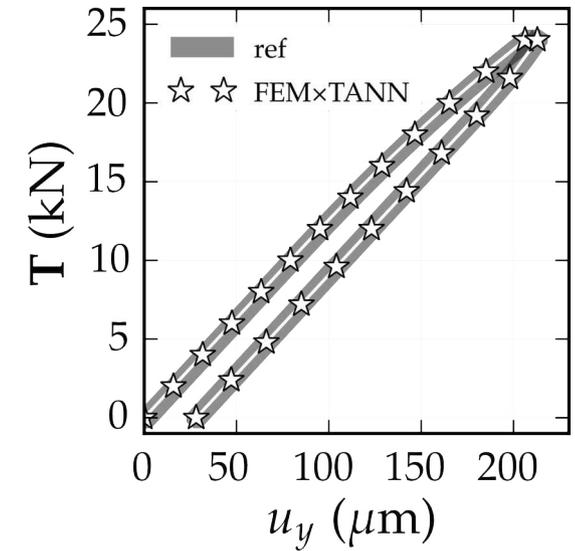
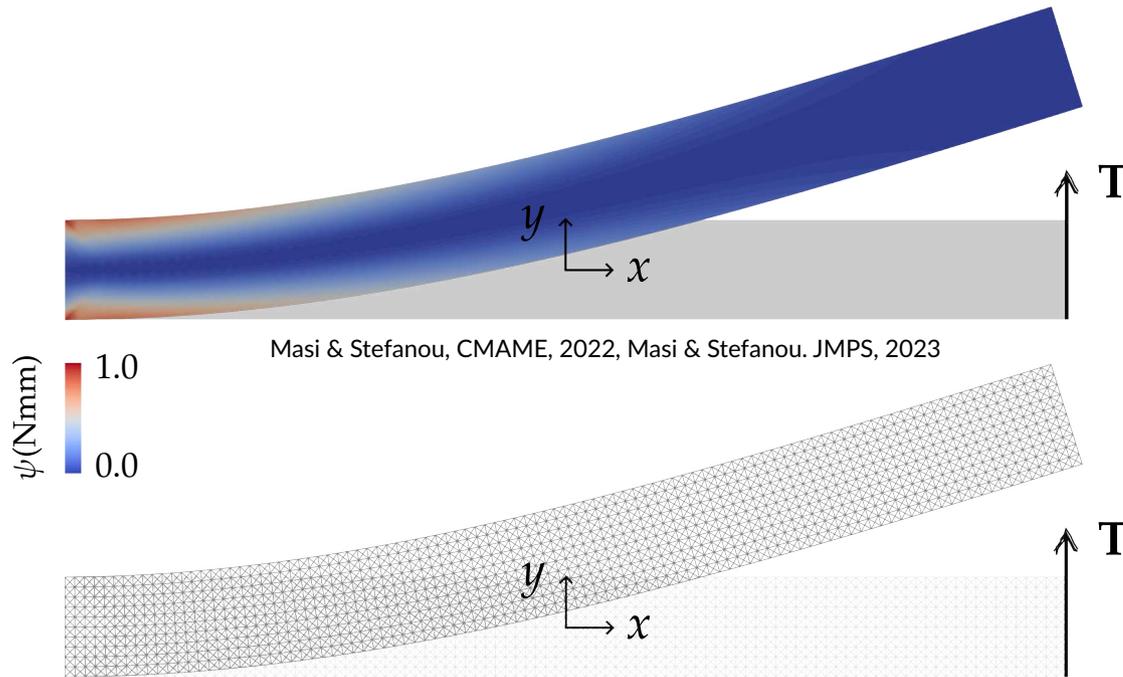


▷ Convergence (asymptotic homogenisation)

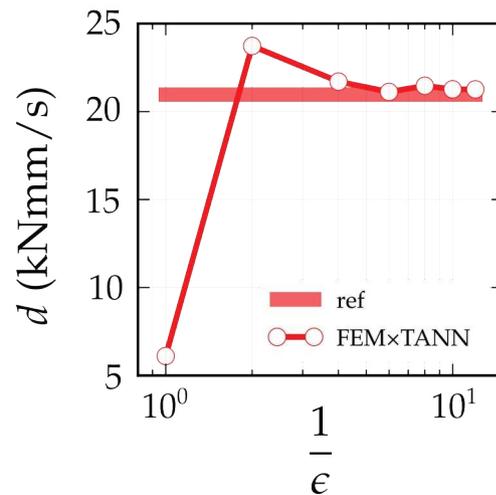
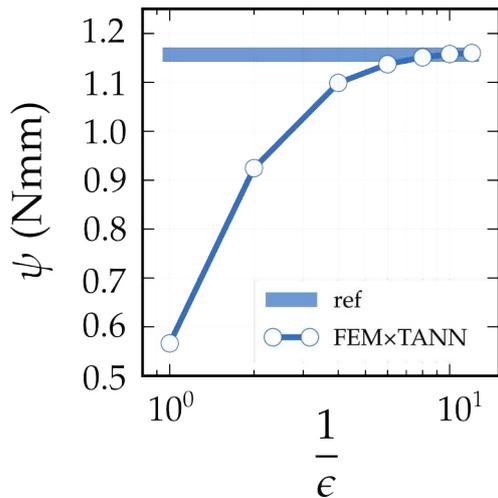


FEM×TANN and numerical homogenisation

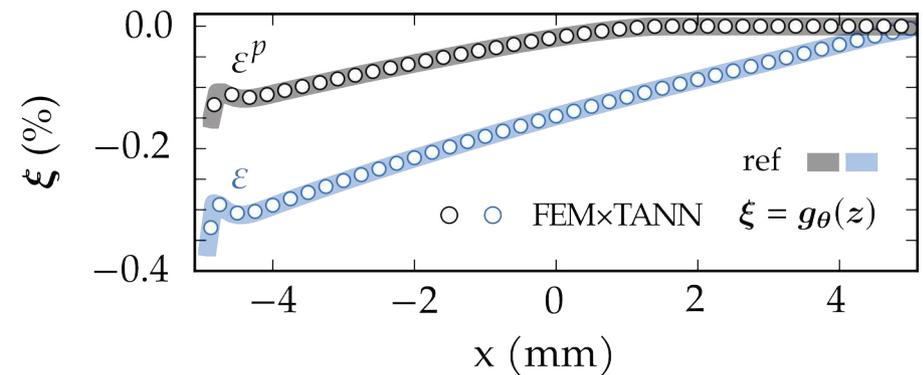
Double-scale homogenisation scheme: rigorous (asymptotic) and fast



▷ Convergence (asymptotic homogenisation)



▷ Microstructure tracking (“localisation”)



I) Background setting

- a) Non-equilibrium thermodynamics

II) Thermodynamics-based artificial neural networks

- a) Methods
- b) Discovery of internal state variables
- c) Experiments

III) Multiscale computing

IV) Neural equations from small data

- a) Data scarcity and sparsity
- b) Methods
- c) Experiments

Data-driven identification from experiments

Big data (measure everything)

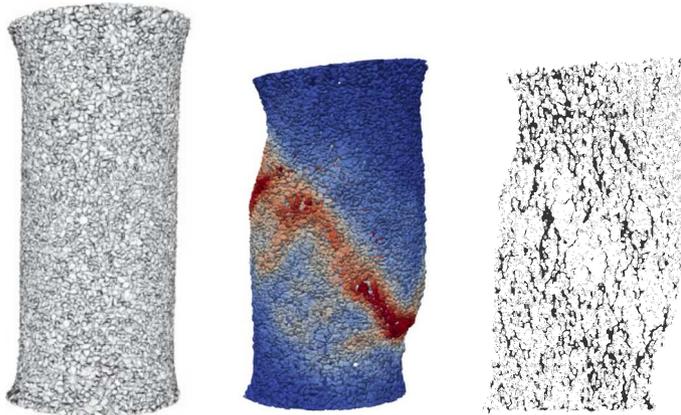
$$\mathcal{D} = \left\{ \boldsymbol{\sigma}^{(n)}, \boldsymbol{\varepsilon}^{(n)}, \mathbb{X}^{(n)} \right\}_{n=0}^N$$

$$= \begin{bmatrix} \boldsymbol{\sigma}^{(0)} & \boldsymbol{\sigma}^{(1)} & \dots & \boldsymbol{\sigma}^{(N)} \\ \boldsymbol{\varepsilon}^{(0)} & \boldsymbol{\varepsilon}^{(1)} & \dots & \boldsymbol{\varepsilon}^{(N)} \\ \mathbb{X}^{(0)} & \mathbb{X}^{(1)} & \dots & \mathbb{X}^{(N)} \end{bmatrix}$$



observe the entire state
(at smaller scales)

all states $\left\{ \mathbb{X}^{(n)} \right\}_{n=0}^N$



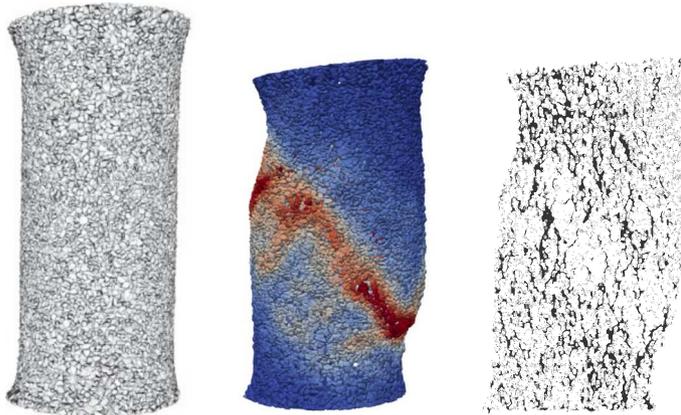
Big data (measure everything)

$$\mathcal{D} = \left\{ \boldsymbol{\sigma}^{(n)}, \boldsymbol{\varepsilon}^{(n)}, \mathbb{X}^{(n)} \right\}_{n=0}^N$$

$$= \begin{bmatrix} \boldsymbol{\sigma}^{(0)} & \boldsymbol{\sigma}^{(1)} & \dots & \boldsymbol{\sigma}^{(N)} \\ \boldsymbol{\varepsilon}^{(0)} & \boldsymbol{\varepsilon}^{(1)} & \dots & \boldsymbol{\varepsilon}^{(N)} \\ \mathbb{X}^{(0)} & \mathbb{X}^{(1)} & \dots & \mathbb{X}^{(N)} \end{bmatrix}$$

observe the entire state
(at smaller scales)

all states $\left\{ \mathbb{X}^{(n)} \right\}_{n=0}^N$



Kawamoto et al. 2018, Karapiperis et al. 2021

Small data (measure observable)

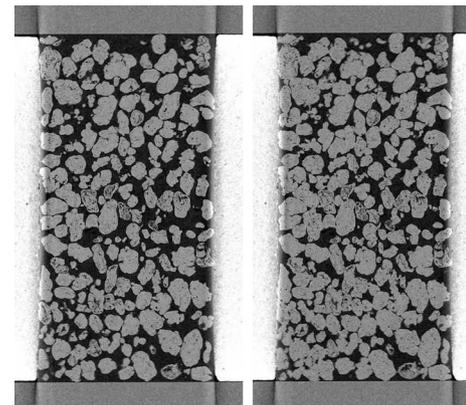
$$\mathcal{D} = \left\{ \boldsymbol{\sigma}^{(n)}, \boldsymbol{\varepsilon}^{(n)} \right\}_{n=0}^N \cup \left\{ \mathbf{x}^{(m)} \right\}_{m=0}^M, \quad M \ll N$$

$$= \begin{bmatrix} \boldsymbol{\sigma}^{(0)} & \boldsymbol{\sigma}^{(1)} & \dots & \boldsymbol{\sigma}^{(N)} \\ \boldsymbol{\varepsilon}^{(0)} & \boldsymbol{\varepsilon}^{(1)} & \dots & \boldsymbol{\varepsilon}^{(N)} \\ \mathbf{x}^{(0)} & ? & ? & \mathbf{x}^{(M)} \end{bmatrix}$$

partial $\mathbf{x} \subset \mathbb{X}$
no observations

$\mathbf{x}^{(0)}$

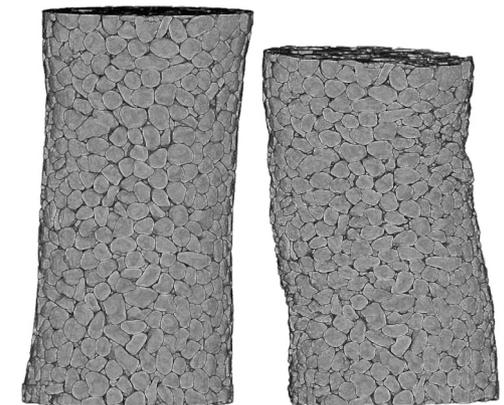
$\mathbf{x}^{(M)}$



Vego, 2023

$\mathbf{x}^{(0)}$

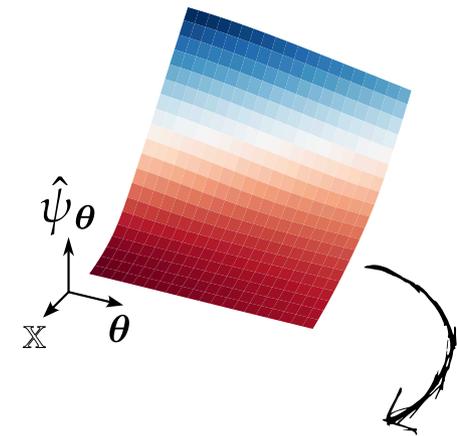
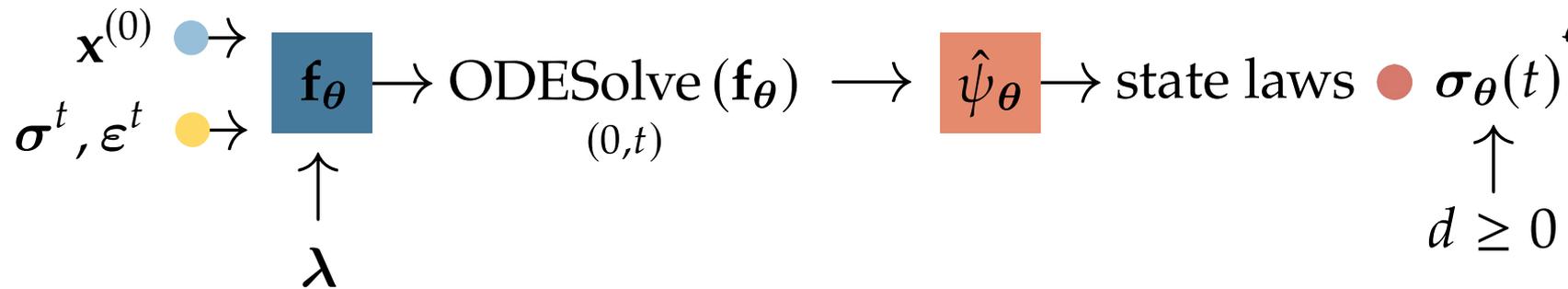
$\mathbf{x}^{(M)}$



Hurley et al. 2023

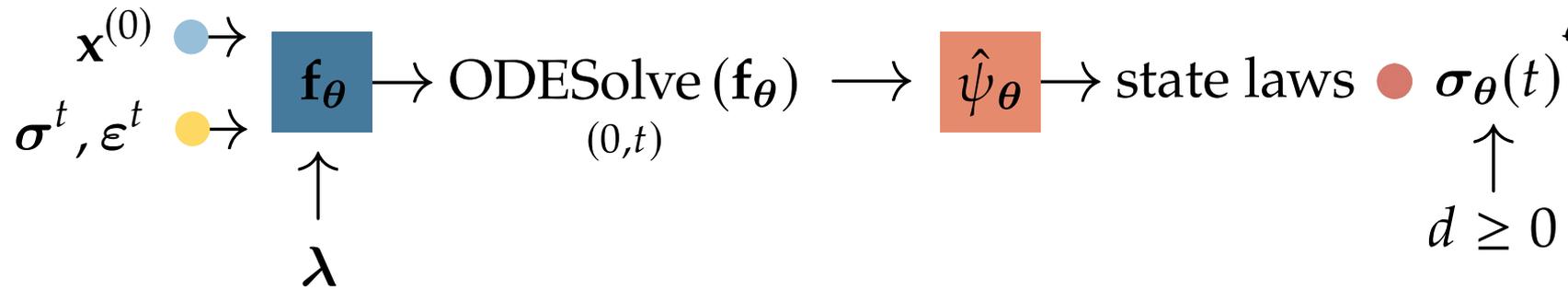
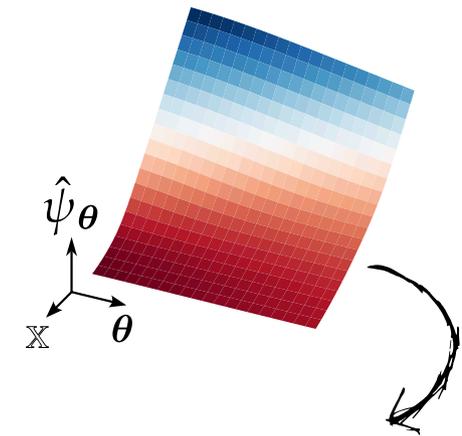
Solving the big problem of small data

Neural Integration for Constitutive Equations (NICE)



Solving the big problem of small data

Neural Integration for Constitutive Equations (NICE)



Optimisation problem: $\theta^*, \lambda^* = \operatorname{argmin} \left(\varrho + \mathcal{L}_\sigma + \mathcal{L}_x + \langle -d_\theta \rangle + \|\theta\|_2^2 \right)$

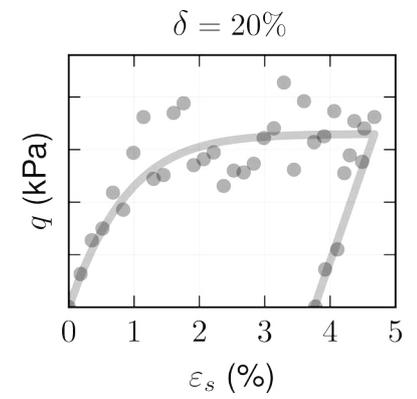
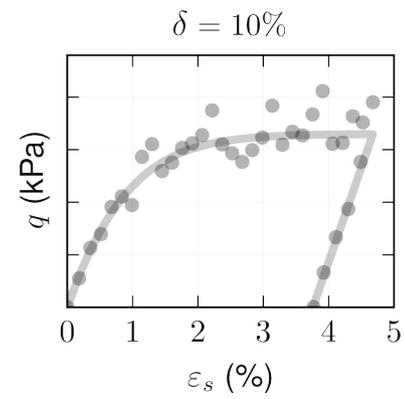
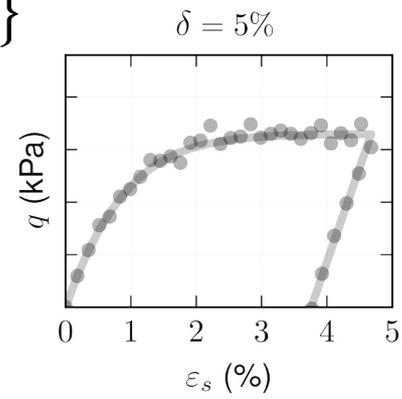
- Residual for initial conditions (elastic strain, etc.) $\varrho = \ell \left(\sigma_\theta \left(\lambda, \mathbf{x}^{(0)} \right) - \sigma^{(0)} \right)$
- Thermodynamics-based loss (stress) $\mathcal{L}_\sigma = \ell \left(\sigma_\theta(t^n), \sigma^{(n)} \right)_{\forall n \in [0, N]}$
- Thermodynamics-based loss (evolution equations) $\mathcal{L}_x = \ell \left(\text{ODESolve}(\mathbf{f}_\theta - \mathbf{x}^{(m)}) \right)_{\forall m \in [1, M]}$

Benchmark

Elasto-plastic material: $\mathbb{X} = \{\epsilon^e\}$

$$p = 1 : \sigma / 3$$

$$q = \sqrt{\frac{3}{2}} \sigma' : \sigma'$$

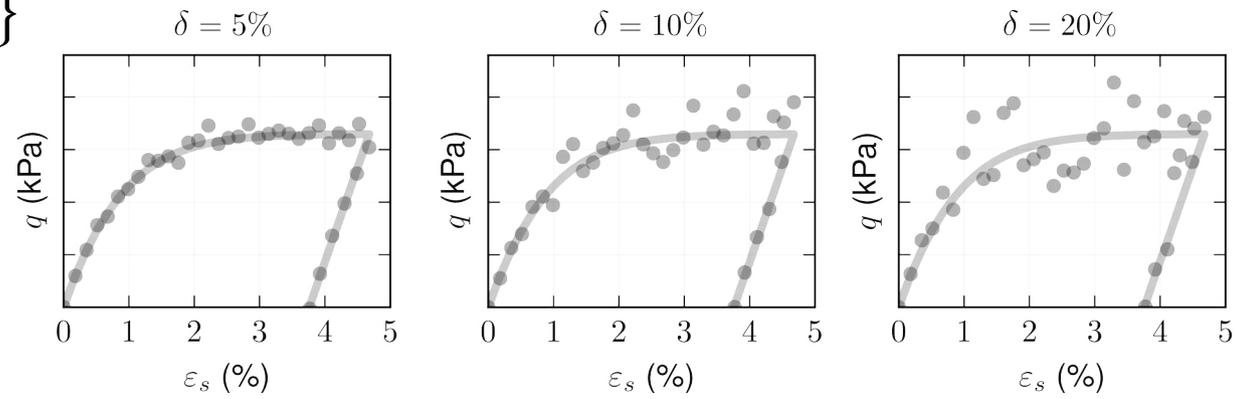


Benchmark

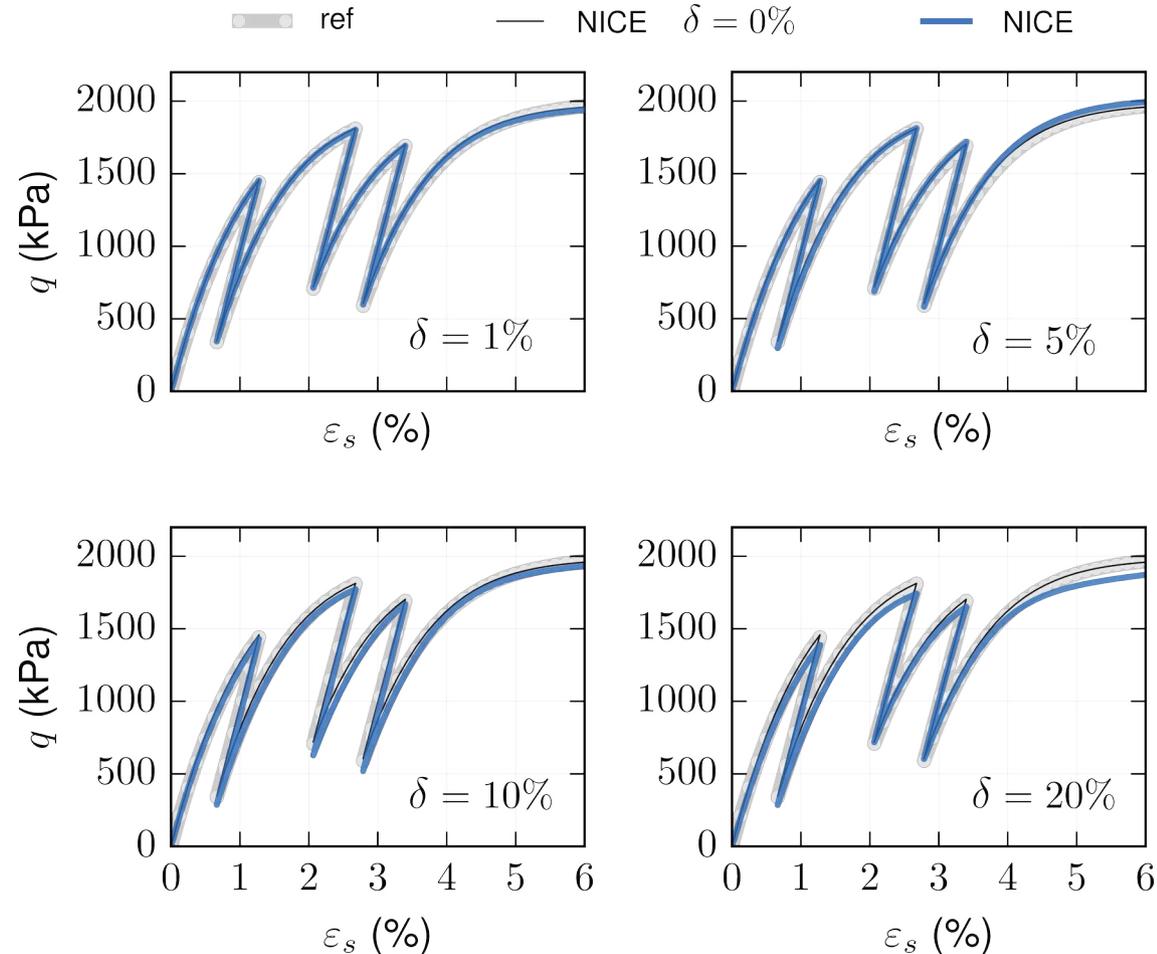
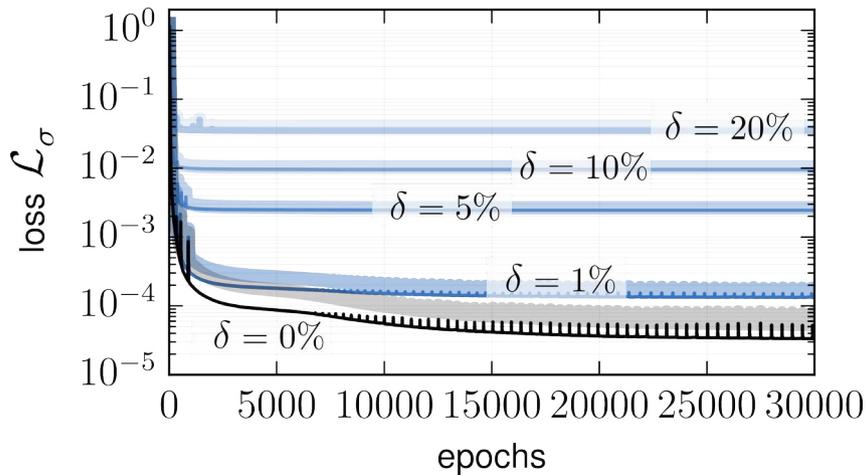
Elasto-plastic material: $\mathbb{X} = \{\epsilon^e\}$

$$p = 1 : \sigma / 3$$

$$q = \sqrt{\frac{3}{2}} \sigma' : \sigma'$$



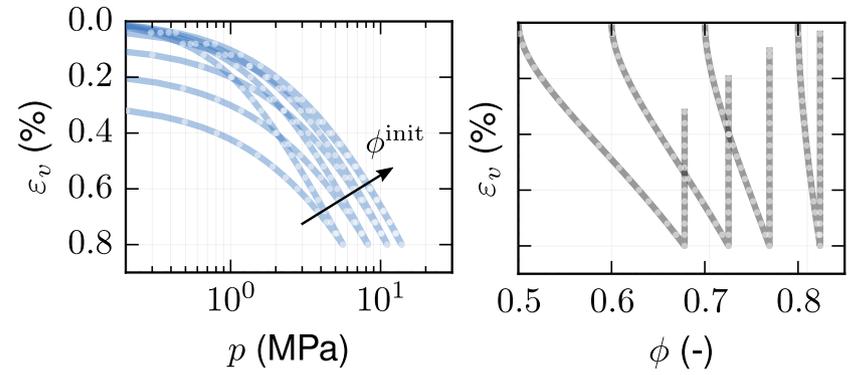
Inference:
unobserved loading path/protocol



▷ High-level generalisation, robustness

Numerical experiment

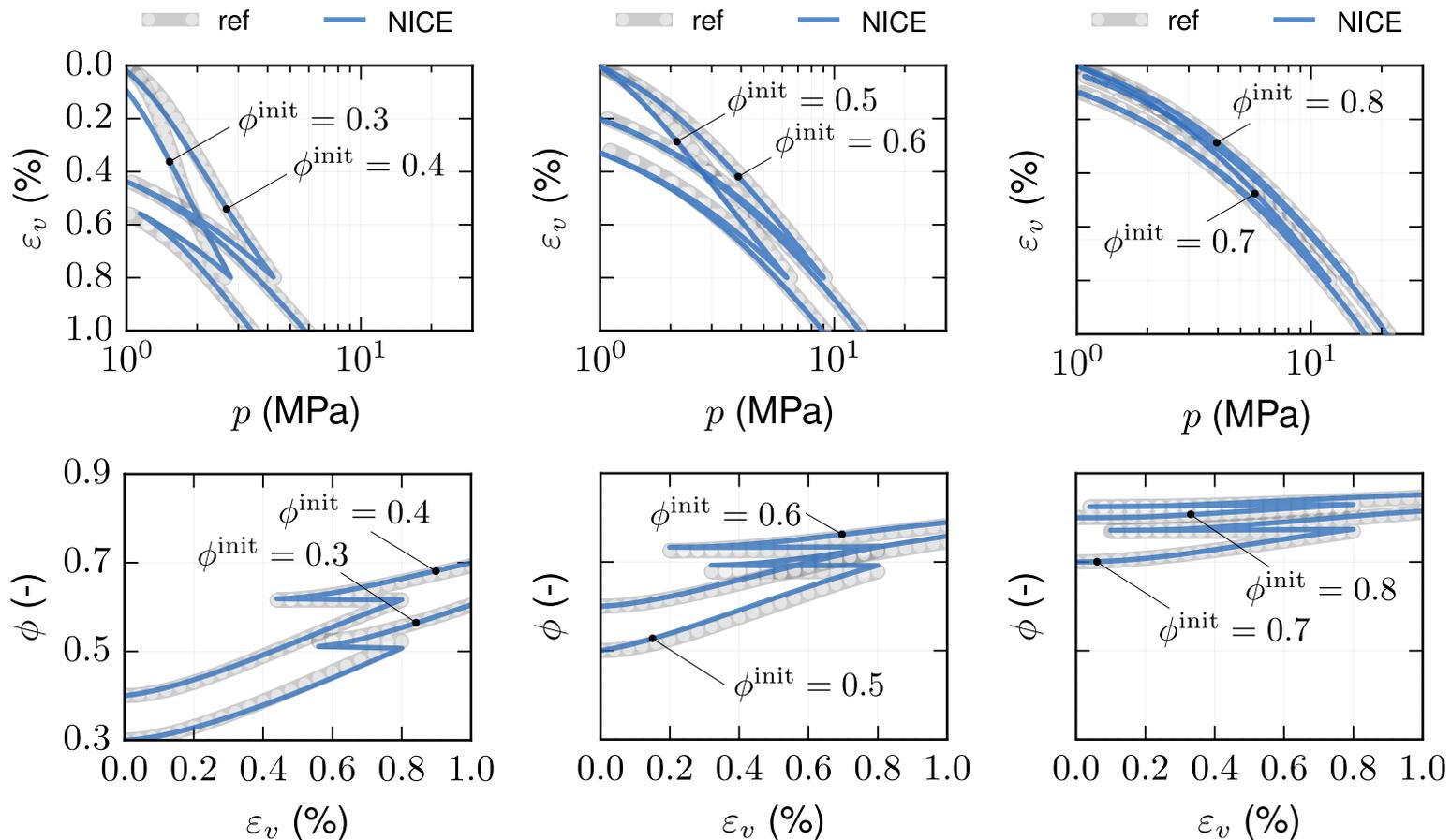
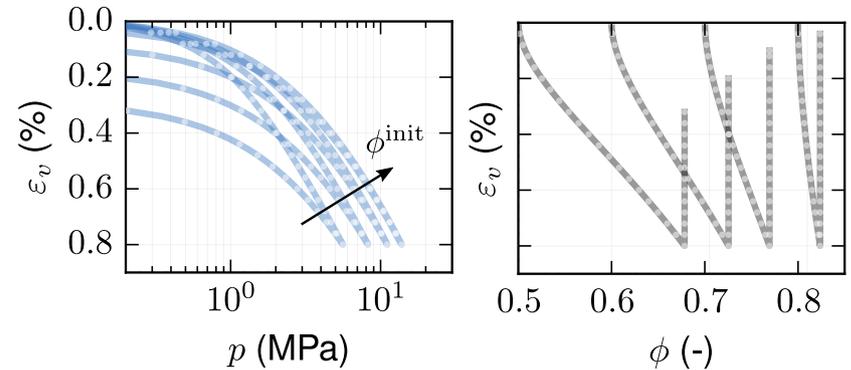
Porous elasto-plastic material: $\mathbb{X} = \{\rho, \varepsilon^e, \mathbf{z}\}$



Numerical experiment

Porous elasto-plastic material: $\mathbb{X} = \{\rho, \varepsilon^e, \mathbf{z}\}$

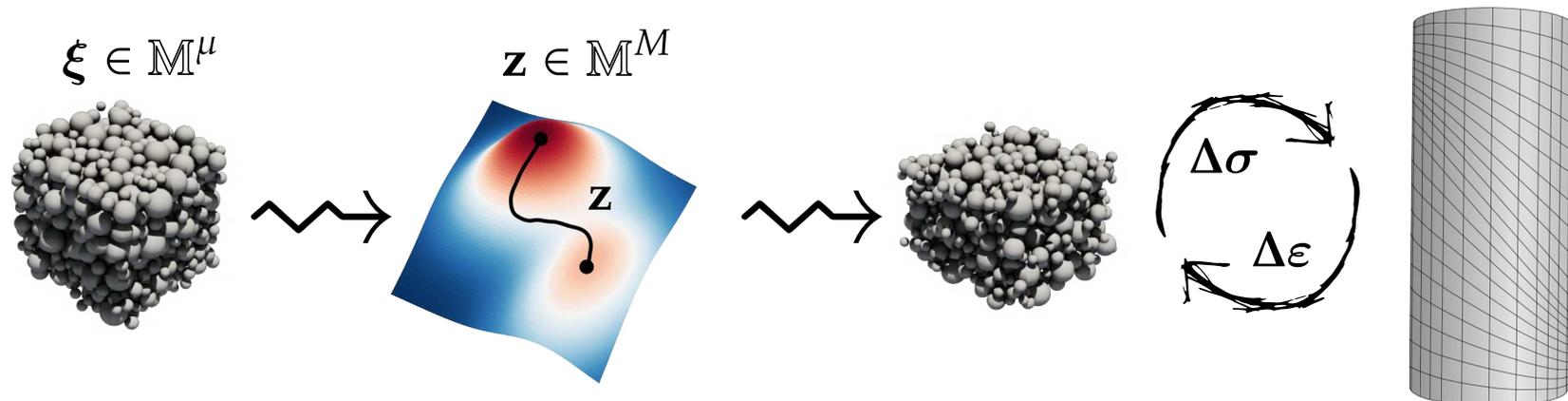
Inference:
unobserved loading path/protocol



▷ High-level generalisation, robustness: discover constitutive equations

Summary

- **Numerical homogenisation** driven by machine learning methods
- Thermodynamics- and data-driven **scale bridging**
- Data from various provenance
(material-point and numerical experiments)
- **Massive speed-ups** of multiscale simulations ($\times 10^3$)
- **History dependence** (plasticity, damage, viscosity, etc.)
- **Discovery of internal descriptors**
- **Learn from small data** (sparse and scarce)



References

- F. Masi et al. (2021). J Mech Phys Solids 147, 104277.
- F. Masi, I. Stefanou (2022). Comput Methods Appl Mech Eng 398, 115190.
- F. Masi, I. Stefanou (2023). J Mech Phys Solids 174, 105245.
- F. Masi (2024). Machine Learning in Geomechanics, vol I. Wiley.
- F. Masi, I. Stefanou (2024). Machine Learning in Geomechanics, vol II. Wiley.
- F. Masi, I. Einav (2024). Comput Methods Appl Mech Eng 420, 116698.

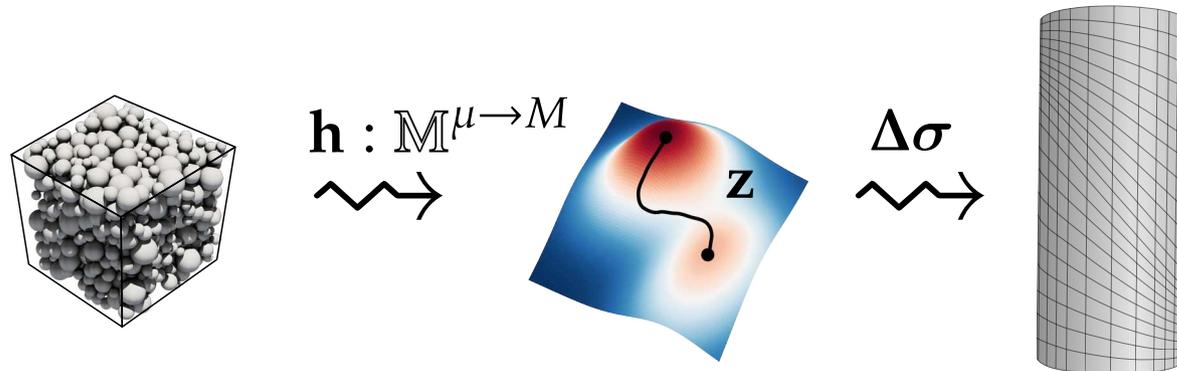
Codes & hands-on

- TANN: github.com/filippo-masi/Thermodynamics-Neural-Networks
- TANN-multiscale: github.com/filippo-masi/TANN-multiscale
- Numerical GeoLab: github.com/AlexSTA1993/numerical_geolab
- NICE: github.com/filippo-masi/NICE
- Hands-on Regression: github.com/alert-geomaterials/2023-doctoral-school
- Hands-on TANN/PINN: github.com/alert-geomaterials/2023-doctoral-school

Acknowledgements



Thermodynamics- and data-driven scale bridging

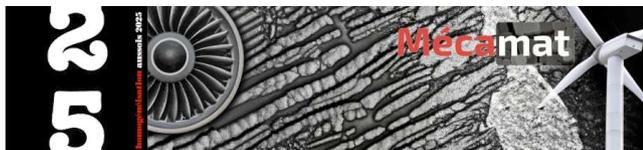


Filippo Masi

Univ. Grenoble Alpes, Inria, CNRS, Grenoble INP, LJK

Mécamat: Homogénéisation du comportement mécanique des matériaux hétérogènes

Aussois – Jan. 27-31, 2025



Inria



UGA
Université
Grenoble Alpes